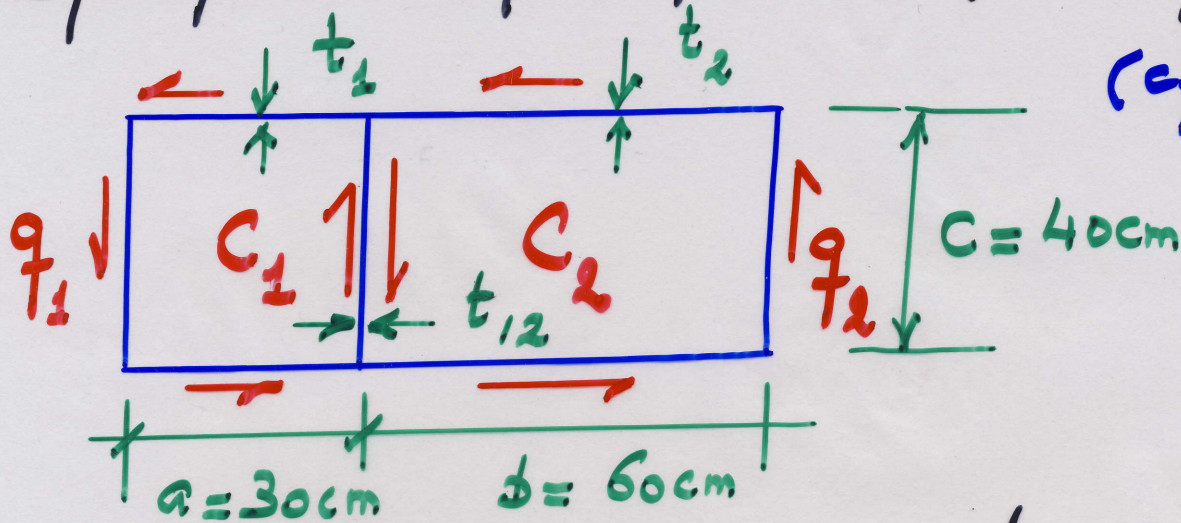


Mtg 19: Wed, 8 Oct 08. EAS 4200C (19-1)

Multicell airfoil: (Sec 3.6)

Specific example first (generalization later)
(cf. p. 94)



$$t_1 = 0.3 \text{ cm}, \quad t_2 = 0.5 \text{ cm}, \quad t_{12} = 0.4 \text{ cm}$$

p. 17-1: Find θ as function of T ,
and J (torsional const.)

$$(1) \quad T = T_1 + T_2 = 2q_1 \bar{A}_1 + 2q_2 \bar{A}_2$$

$$\bar{A}_1 = ac, \quad \bar{A}_2 = bc$$

$$\theta_1 = \frac{1}{2G\bar{A}_1} \oint \frac{q_1}{t_1} ds$$

$$(3) \quad = \frac{1}{2G\bar{A}_1} \left[\frac{2q_1 a}{t_1} + \frac{q_1 c}{t_1} + \frac{(q_1 - q_2)c}{t_{12}} \right]$$

$$\theta_2 = \frac{1}{2G \bar{A}_2} \oint \frac{q_2}{t_2(s)} ds \quad \underline{19-2}$$

$$\stackrel{(4)}{=} \frac{1}{2G \bar{A}_2} \left[\frac{2q_2 b}{t_2} + \frac{q_2 c}{t_2} + \frac{(q_2 - q_1)c}{t_{12}} \right]$$

C_1 and C_2 have same rate of twist angle: $\theta_1 = \theta_2$ (2) constraint eq. on θ_i 's.

Think of (q_1, q_2) as 2 unknowns. Eqs (1) & (2) are 2 eqs. for the 2 unknowns (q_1, q_2) w/ T being a variable. \Rightarrow use (1) & (2) to find expr. for (q_1, q_2) in terms of T .

$$\begin{cases} q_1 = \beta_1 T \\ q_2 = \beta_2 T \end{cases} \quad | \quad (\beta_1, \beta_2) \text{ are actually numbers.} \quad \underline{HW.}$$


Next, use the expr. (3) or (4) (19-3)
to find the expr. betw. θ and T :

$$\theta = \theta_1 = \theta_2 = \frac{T}{2GJ}$$

and deduce J .

Recall: $T = GJ\theta$

Once θ is found as a func. of
 T , use $\theta = \frac{T}{GJ}$ or $J = \frac{T}{G\theta}$
to find J .



Mtg 20: Fri, 10 Oct 08. EAS 4200c (20-1)

HW 4: 3-cell NACA 2415 air foil.

p. 16.2: 2 partition walls

1st wall at $\frac{1}{4} c$ from leading edge

2nd " " $\frac{3}{4} c$ " " "

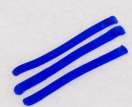
3 unknowns q_1, q_2, q_3


$$T = 2 \sum_{i=1}^3 q_i \bar{A}_i \quad (1)$$

$$\theta_1 = \theta_2 \quad (2)$$

$$\theta_2 = \theta_3 \quad (3)$$

Find J .

Note: need thickness t , 

but G cancels out. 

Now theory (deriv.)

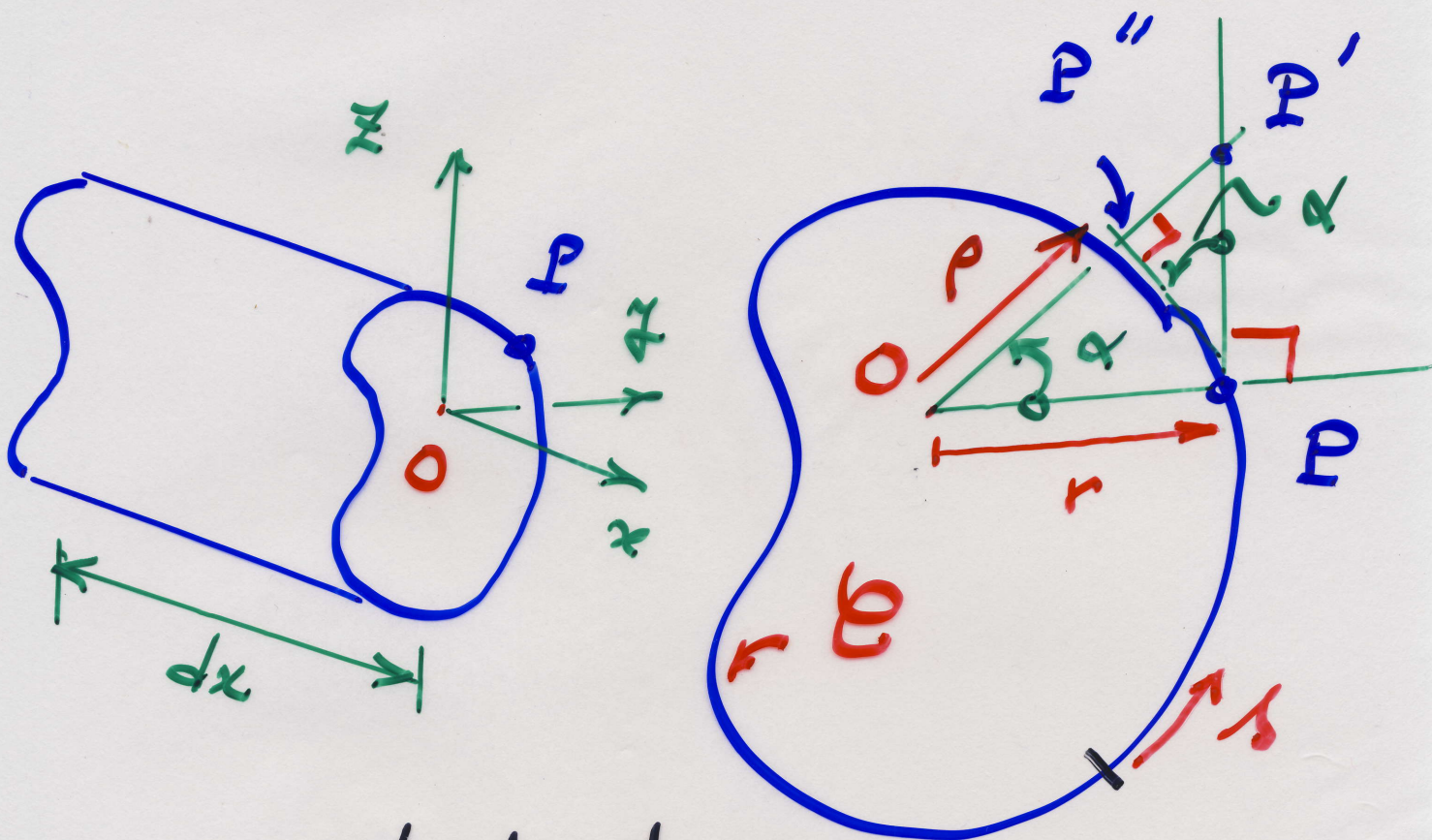
$$T = G J \theta \quad (\text{done})$$

$$T = 2 q \bar{A} \quad (\text{done})$$

Engineering (ad-hoc) deriv. of Eq-2

$$\theta = \frac{1}{2GA} \oint \frac{q}{t} ds$$

uniform bar w/ non-circular cross section subject. twist:



Disp PP' due to α :

$$\frac{PP'}{OP} = \tan \alpha \approx \alpha \quad (\text{for } \alpha \text{ small})$$

Proj. disp PP'' on the dir. $\perp OP'$:

$$PP'' = PP' \cos \alpha$$

$$PP'' = (OP \tan \alpha) \cos \alpha \quad (20-3)$$

$$= \underbrace{(OP \cdot \cos \alpha)}_{OP''} \tan \alpha$$

Recall: $\begin{cases} OP = r & (\text{radial coord.}) \\ OP'' = \rho & (\text{p. 12-5}) \end{cases}$

$$PP'' = \underbrace{(r \cos \alpha)}_{\rho} \underbrace{\tan \alpha}_{\alpha}$$

↓
 disp. of P
 in dir. "tangent" to lateral surf.
 of bar.

Strain: $\gamma = \frac{PP''}{dx} = \frac{\rho \alpha}{dx} = \rho \theta$

with $\theta = \frac{\alpha}{dx}$ rate of twist

Rafael. $\left\{ \begin{array}{l} (\alpha \text{ very small}) \\ \text{can be denoted} \\ \text{by } d\alpha \end{array} \right.$

Mtg 21: Mon, 13 Oct 08. EAS 4200C (21-1)

$$\theta = \frac{d\alpha}{dx}$$

Hooke's law: $\tau = G\gamma = G\rho\theta$

$$\tau(s) = G\rho(s)\theta(x)$$

Int. along contour \mathcal{C} :

$$\int_{\mathcal{C}} \underbrace{\tau(s)}_{\substack{q(s) \\ t(s)}} ds = G\theta(x) \int_{\mathcal{C}} \underbrace{\rho(s)}_{2\bar{A}} ds$$

Hence expr. for θ p. 20-2.

Q: What is ad-hoc about the above deriv. of θ expr. on p. 20-2 and the deriv. of $T = 2q\bar{A}$?

- 1) Strain γ must be obtained using the disp. of P in the dir. tangent to \mathcal{C} at P , but PP''

on γ . 20-2 is not nec. tangent $\underline{L}21-2$
to \mathcal{C} (but actually close).

2) $\tau = \frac{q}{t}$ obtained from ad hoc
assump. that τ was uniform
across wall thickness.

Now formal justification (deriv.)
by elasticity theory:

Roadmap γ . 16-2

A. Kinematic assump. γ . 16-1.

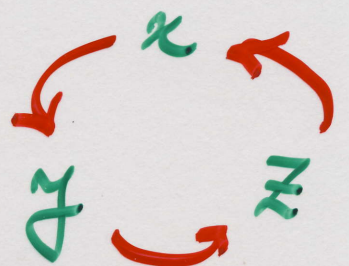
$$u_x(\gamma, z) = \theta \psi(\gamma, z)$$

\uparrow
considered const.
wrt x : uniform bar

$$u_\gamma(x, z) = -\theta x z$$

$$u_z(x, \gamma) = \theta x \gamma$$

To transf. eqs. in Sun [2006] (21-3)
 to those using our unified notation,
 use cyclic permutation



(1) $\epsilon_{xx} = \epsilon_{yy} = \epsilon_{zz} = \gamma_{yz} = 0$

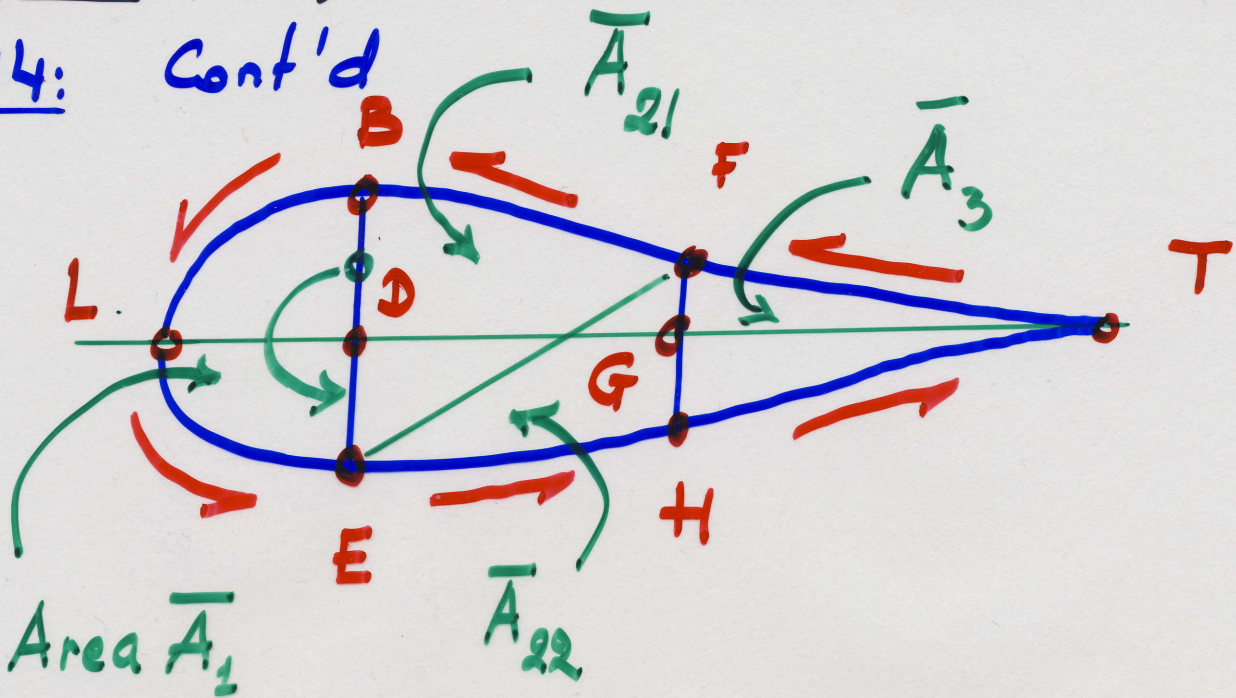
$$\epsilon_{xx} = \frac{\partial u_x(x, z)}{\partial x} = 0$$

$$\gamma_{yz} = \underbrace{\frac{\partial u_y(x, z)}{\partial z}}_{-\theta_x} + \underbrace{\frac{\partial u_z(x, y)}{\partial y}}_{+\theta_x} = 0$$

HW: Do ϵ_{yy} , ϵ_{zz} .

Mtg 22: Wed, 15 Oct 08. EAS 4200C (22-1)

HW4: Cont'd



1) Set $P_0 = D$, sweep area $B \rightarrow L$
 $\rightarrow E$

Area \bar{A}_{21}

2) Set $P_0 = E$, sweep area $F \rightarrow B$

2.1

2.2) Area \bar{A}_{22} : Set $P_0 = F$, sweep
area $E \rightarrow H$

3) Area \bar{A}_3 : Set $P_0 = G$, sweep area
 $H \rightarrow T \rightarrow F$

$$\bar{A} = \bar{A}_1 + \underbrace{\bar{A}_{21} + \bar{A}_{22}}_{\bar{A}_2} + \bar{A}_3$$

within 1% of previous results

\bar{A}_2



p. 21-2 cont'd : 3rd "ad-hoc" pt 22-2
in engineering deriv.

3) Inconsistency in assump. on size of

p. 20-2

α : To get line PP' , assume α small ; to get PP'' , assume α finite ($\cos \alpha$) and $p = r \cos \alpha$; reintroduce small α after that.

Back to formal deriv. (p. 21-3)

How many strain comp. in 3-D : 6

$$\underline{\underline{\epsilon}} = \begin{bmatrix} \epsilon_{xx} & \epsilon_{xy} & \epsilon_{xz} \\ \epsilon_{yx} & \epsilon_{yy} & \epsilon_{yz} \\ \epsilon_{zx} & \epsilon_{zy} & \epsilon_{zz} \end{bmatrix}$$

9 coeff
(Diego)

3x3

$x \leftrightarrow 1$
 $y \leftrightarrow 2$
 $z \leftrightarrow 3$ } indicial notation

$$\underline{\underline{\epsilon}} = \begin{bmatrix} \epsilon_{11} & \epsilon_{12} & \epsilon_{13} \\ \epsilon_{21} & \epsilon_{22} & \epsilon_{23} \\ \epsilon_{31} & \epsilon_{32} & \epsilon_{33} \end{bmatrix} = [\epsilon_{ij}]$$

$i, j = 1, 2, 3$

tensorial

ϵ_{ij}
row \nearrow \nwarrow col.

Sym. of $\underline{\underline{\epsilon}}$:

$$\epsilon_{ij} = \frac{1}{2} \left(\frac{\partial u_i}{\partial x_j} + \frac{\partial u_j}{\partial x_i} \right)$$

Coord. $x \leftrightarrow x_1$
 $y \leftrightarrow x_2$
 $z \leftrightarrow x_3$

$$\epsilon_{11} = \epsilon_{xx} = \frac{1}{2} \left(\frac{\partial u_1}{\partial x_1} + \frac{\partial u_1}{\partial x_1} \right)$$

$$= \frac{\partial u_1}{\partial x_1} = \frac{\partial u_x}{\partial x}$$

Sym: $\epsilon_{ij} = \epsilon_{ji}$ (HW)

$$\epsilon_{12} = \epsilon_{21} \iff \epsilon_{xy} = \epsilon_{yx} \quad (22-4)$$

Hence: only 6 indep. comp. of $\underline{\epsilon}$ in 3D
Similarly for stress tensor $\underline{\sigma}$

$$\underline{\sigma} = [\sigma_{ij}]_{3 \times 3}$$

Also only 6 indep. comp. of $\underline{\sigma}$ in 3-D

Q: (J_{eff}) Does sym. of $\underline{\epsilon}$ related to isotropy of materials?

No: Isotropic elasticity is related to material behavior ($\underline{\sigma}$ - $\underline{\epsilon}$ rel.)

Following Eq. (1) on p. 21-3, ≡

$$\sigma_{xx} = \sigma_{yy} = \sigma_{zz} = \tau_{yz} = 0$$

due to σ - ϵ rel. why? (cf. p. 70)

Consider isotropic elastic mat'l:

\uparrow mat'l behavior

Mtg 23: Fri, 17 Oct 08. EAS 4200C (23-1)

Read: MIT OCW Prof. Lagace Unit 4,
How to relate strains to E, ν ?

Stere: Hooke's law.

↑ ↑ "nu"
Young's modulus Poisson's ratio

Normal strains:

$$\epsilon_{xx} = \frac{\sigma_{xx}}{E} - \nu \frac{\sigma_{yy}}{E} - \nu \frac{\sigma_{zz}}{E}$$

$$\epsilon_{yy} = HW$$

$$\epsilon_{zz} = HW$$

Shear strains:

$$\tau_{xy} = 2\epsilon_{xy} = \frac{\tau_{xy}}{G}$$

$\tau_{xy} = \sigma_{xy}$
etc.

$$\tau_{yz} = HW$$

$$\tau_{zx} = HW$$

Voigt notation:

Due sym. of $\underline{\epsilon} = [\epsilon_{ij}]$, $\underline{\sigma} = [\sigma_{ij}]$
arrange ϵ_{ij} 's and σ_{ij} 's in col. mat.

$$\{\epsilon_{ij}\}_{6 \times 1} = \begin{Bmatrix} \epsilon_{11} \\ \epsilon_{22} \\ \epsilon_{33} \\ \epsilon_{23} \\ \epsilon_{31} = \epsilon_{13} \\ \epsilon_{12} \end{Bmatrix} \quad \{\sigma_{ij}\}_{6 \times 1} = \begin{Bmatrix} \sigma_{11} \\ \sigma_{22} \\ \sigma_{33} \\ \sigma_{23} \\ \sigma_{31} = \sigma_{13} \\ \sigma_{12} \end{Bmatrix}$$

Hooke's law for isotropic elasticity:

$$\begin{Bmatrix} \epsilon_{11} \\ \epsilon_{22} \\ \vdots \\ \gamma_{yz} \\ \gamma_{xz} \\ \gamma_{xy} \\ \epsilon_{12} \end{Bmatrix} = \begin{bmatrix} \frac{1}{E} & -\frac{\nu}{E} & -\frac{\nu}{E} & 0 & 0 & 0 \\ \text{HW} & & & & & \\ \text{HW} & & & & & \\ \hline 0 & 0 & 0 & \frac{1}{G} & 0 & 0 \\ \text{HW} & & & & & \\ \text{HW} & & & & & \end{bmatrix} \begin{Bmatrix} \sigma_{11} \\ \sigma_{22} \\ \vdots \\ \vdots \\ \vdots \\ \sigma_{12} \end{Bmatrix}$$

$\gamma_{yz} = \gamma_{23}$ (Jared, Greg)

HW: Fill out all rows (Chris)

(23-2)

$$\begin{Bmatrix} \epsilon_{11} \\ \epsilon_{22} \\ \epsilon_{33} \\ \epsilon_{23} \\ \epsilon_{31} \\ \epsilon_{12} \end{Bmatrix} = \begin{bmatrix} & & & & & \\ & & & & & \\ & & & & & \\ & & & & & \\ & & & & & \\ & & & & & \\ & & & & & \end{bmatrix} \begin{Bmatrix} \sigma_{11} \\ \sigma_{22} \\ \sigma_{33} \\ \sigma_{23} \\ \sigma_{31} \\ \sigma_{12} \end{Bmatrix}$$

$\frac{1}{2G}$

Note: Steel, $\nu = 0.3$
Cork, $\nu = 0$
rubber, $\nu = 0.5$ (incompressible)
regular mat'l : $0 \leq \nu \leq 0.5$
Artificial neg. Poisson's ratio mat'l
 $\nu < 0$

Mtg 24: Mon, 20 Oct 08. EAS 4200c (24-1)

- p. 21.2: kinematic assump. (")
- p. 21-3: 4 zero strain comp. (cf. p. 70)
- p. 22-4: 4 zero stress comp. (")
- p. 23-2: $\epsilon - \sigma$ rel.
(2 forms: tensorial or engineering)
p. 23-2 p. 23-1

p. 23.2: Rewrite $\epsilon - \sigma$ rel:

$$\{\epsilon_{ij}\}_{6 \times 1} = \underbrace{\begin{bmatrix} \underline{A}_{3 \times 3} & \underline{0}_{3 \times 3} \\ \underline{0}_{3 \times 3} & \underline{B}_{3 \times 3} \end{bmatrix}}_{6 \times 6} \{\sigma_{ij}\}_{6 \times 1}$$

full (not diag.) diag

$\sigma - \epsilon$ rel:

$$\{\sigma_{ij}\}_{6 \times 1} = \underbrace{\begin{bmatrix} \underline{A}^{-1}_{3 \times 3} & \underline{0}_{3 \times 3} \\ \underline{0}_{3 \times 3} & \underline{B}^{-1}_{3 \times 3} \end{bmatrix}}_{6 \times 6} \{\epsilon_{ij}\}_{6 \times 1}$$

full diag Id. mat.

Verification: (Adam)

$$\underline{C}^{-1}_{6 \times 6} \underline{C}_{6 \times 6} = \underline{I}_{6 \times 6}$$

$$C^{-1} C = \begin{bmatrix} \underline{A^{-1}} \underline{A} & \underline{0} \\ \underline{0} & \underline{B^{-1}} \underline{B} \end{bmatrix} = \underline{I} \quad \text{[24-2]}$$

p. 21-3:

$$\{\sigma_{ij}\} = \begin{bmatrix} \underline{A^{-1}} & \underline{0} \\ \underline{0} & \underline{B^{-1}} \end{bmatrix} \left\{ \begin{array}{c} 0 \\ 0 \\ 0 \\ \hline 0 \\ \epsilon_{31} \\ \epsilon_{12} \end{array} \right\}$$

↑
Diag.

$$\sigma_{11} = 0$$

$$\sigma_{22} = 0$$

$$\sigma_{33} = 0$$

$$\sigma_{23} = 2G \epsilon_{23}$$

$$\sigma_{31} = 2G \epsilon_{31}$$

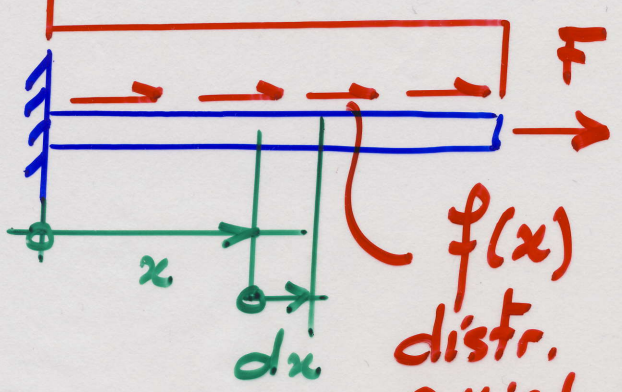
$$\sigma_{12} = HW$$

$$0 = \frac{1}{2} \gamma_{23} = \frac{1}{2} \gamma_{32}$$

Roadmap (p. 16-2): C. Equil. Eq. for stresses
in a non uniform stress field (Sec. 2.4)

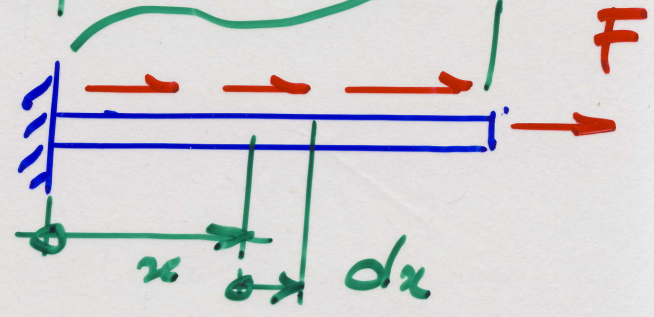
Consider 1-D case as model.

$f(x)$ Uniform (Sec 2.3)



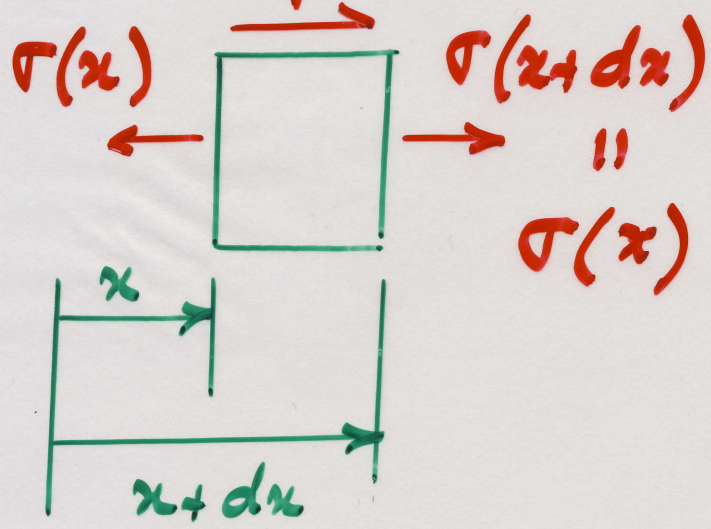
$f(x) = \text{const}$
 $= 0$
 distr. axial load

$f(x)$ Nonuniform

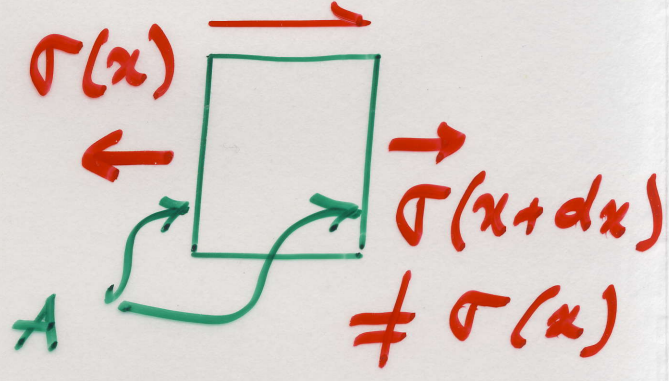


$f(x) = \text{non-const.}$

$f(x) = 0$

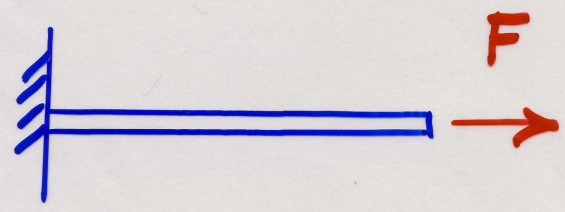


$f(x) \neq 0$



Case includes

$f(x) = \text{const} \neq 0$



Mtg 26: Mon, 27 Oct 08

(26-0)

(Mtg 25: Exam 2 on Wed, 22 Oct 08)

- Biaxial bending (Sec. 4.2)
Directional

- Example 4.1, p. 124

- HWT: Will give areas of stringers A_B, A_E
 A_H, A_F
in e-mail.

- Continue Roadmap for Torsional analysis
* Eq. of equil. for stresses.

Mtg 26: Mon, 27 Oct 08. EAS 4200 C (26-1)

(Mtg 25: Exam 2, Wed, 22 Oct 08)

- Bidirectional Bending (Sec. 4.2), Read 4.1.

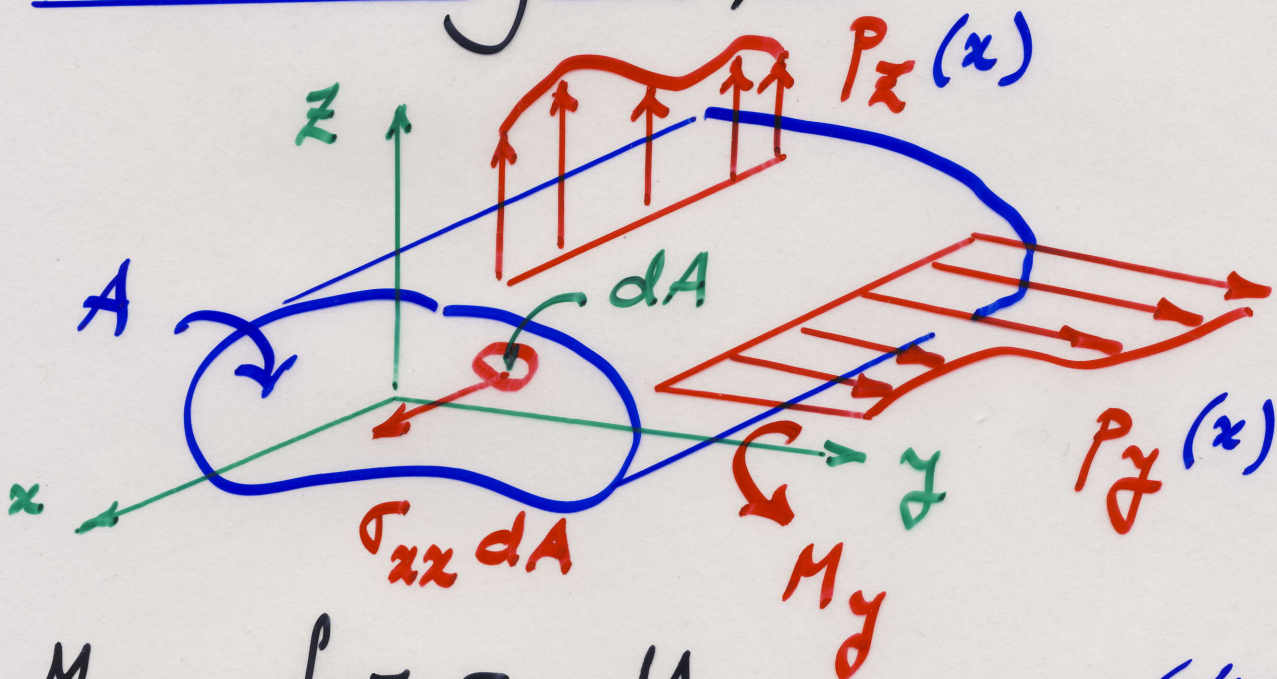
* Example 4.1, p. 124

HWS: Stringers @ pts B, E, H, F (p. 22-1)

Skin and spar web do not participate in bending; only stringers.

Need areas of stringers: A_B, A_E, A_H, A_F

Bidir. bending recipe:



$$M_y = \int_A z \sigma_{xx} dA \quad (4.26)$$

Similarly for M_z (4.27)

Moment of inertia tensor: I_y, I_z, I_{yz} (26.2)

$$I_y = \int_A z^2 dA \quad (4.28a) \quad I_{yy} = I_{zz} = I_{yz}$$

$$I_z = \int_A y^2 dA \quad (4.28b) \quad I_{22} = I_{33}$$

$$I_{yz} = \int_A yz dA \quad (4.28c)$$

$$\begin{aligned} \sigma_{xx} &= E \epsilon_{xx} \\ &= \frac{I_y M_z - I_{yz} M_y}{[I_y I_z - (I_{yz})^2]} y \\ &+ \frac{I_z M_y - I_{yz} M_z}{[I_y I_z - (I_{yz})^2]} z \end{aligned} \quad (4.29)$$

Handwritten annotations: Green arrows point from m_y to the first term and from m_z to the second term. Blue arrows point from D to the denominators.

Recall:

$$\underline{I} = \begin{bmatrix} I_{11} & I_{12} & I_{13} \\ & I_{22} & I_{23} \\ & & I_{33} \end{bmatrix}$$

Sym.

Note: $D := I_y I_z - (I_{yz})^2$ (26-3)

$$= I_{22} I_{33} - (I_{23})^2$$

is the determinant of $\begin{bmatrix} I_{22} & I_{23} \\ I_{32} & I_{33} \end{bmatrix}$
 (Jared)

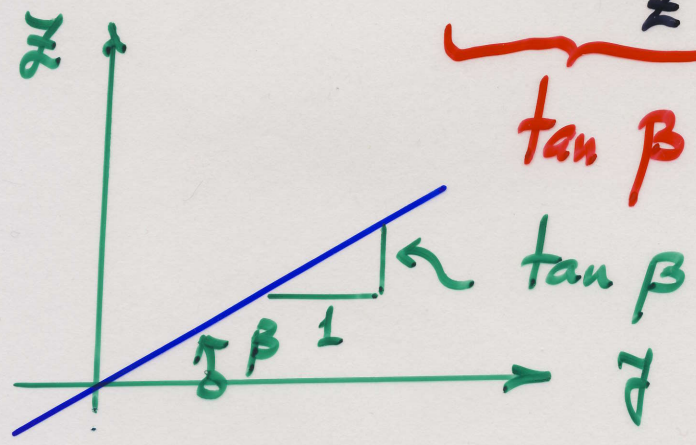
$$I_{32} = I_{zy} \stackrel{\text{def}}{=} \int_A \underbrace{zy}_{\text{yz}} dA = I_{yz} = I_{23}$$

Neutral axis: where $\sigma_{xx} = 0$

$$(4.29) \Rightarrow \sigma_{xx} = m_y \bar{y} + m_z \bar{z} = 0$$

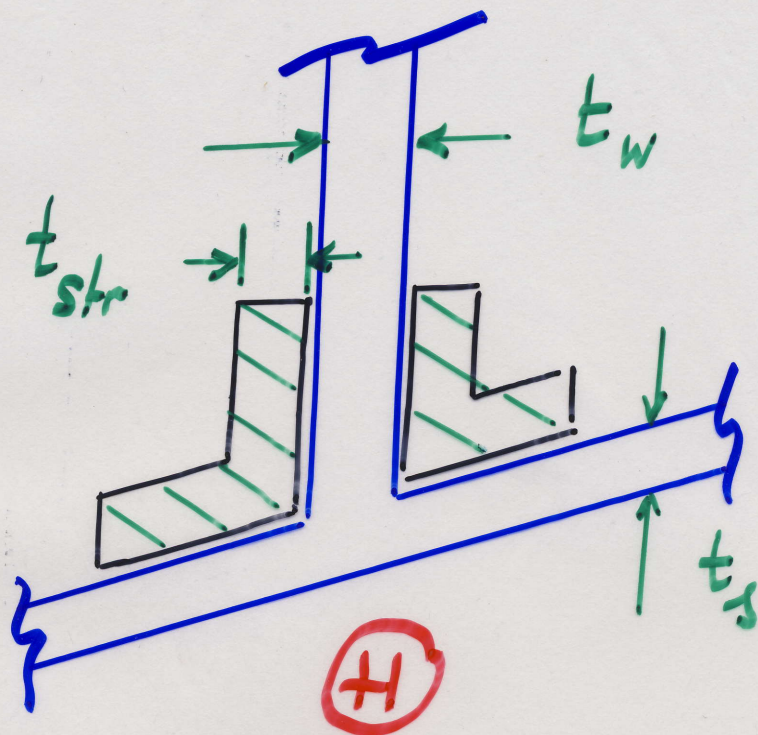
$$\Rightarrow \bar{z} = \left(- \frac{m_y}{m_z} \right) \bar{y} = (\tan \beta) \bar{y}$$

$$\tan \beta = - \tan \alpha \quad (4.30)$$



Wed, 29 Oct 08. EAS 4200C
Mtg 27
HLW 5 Cont'd

27-1



Cont. from p. 24-3 : Eq. of equil. in terms
of $\underline{\sigma}$

Goal:

$$\frac{\partial \sigma_{yx}}{\partial y} + \frac{\partial \sigma_{zx}}{\partial z} = 0$$

cf. (3.14)
(1)

Let's use indicial notation

$$\frac{\partial \sigma_{21}}{\partial x_2} + \frac{\partial \sigma_{31}}{\partial x_3} = 0$$

(2)

Recall: (Sec. 2.4)

(27-2)

$$(3) \quad \frac{\partial \sigma_{xx}}{\partial x} + \frac{\partial \sigma_{yx}}{\partial y} + \frac{\partial \sigma_{zx}}{\partial z} = 0 \quad \text{cf. (2.21)}$$

Similarly, see (2.22) & (2.23) HW5

In indicial notation, (2.21), (2.22), (2.23):

$$(4) \quad \boxed{\sum_{i=1}^3 \frac{\partial \sigma_{ij}}{\partial x_i} = 0 \quad \text{for } j=1, 2, 3}$$