Special relativity and steps towards general relativity: ϵGR

GR: intro

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= vector space (e.g. 4-momentum vectors)

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= dual vector space (think: contour map, gradients)

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duality in a basis of $T_{\mathbf{x}}M$ and a basis of $T^*_{\mathbf{x}}M$ usually defined using δ^μ_ν

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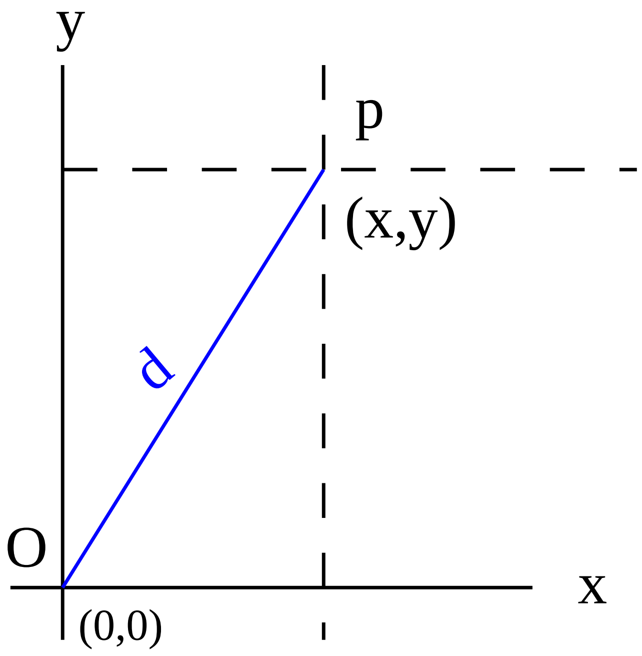
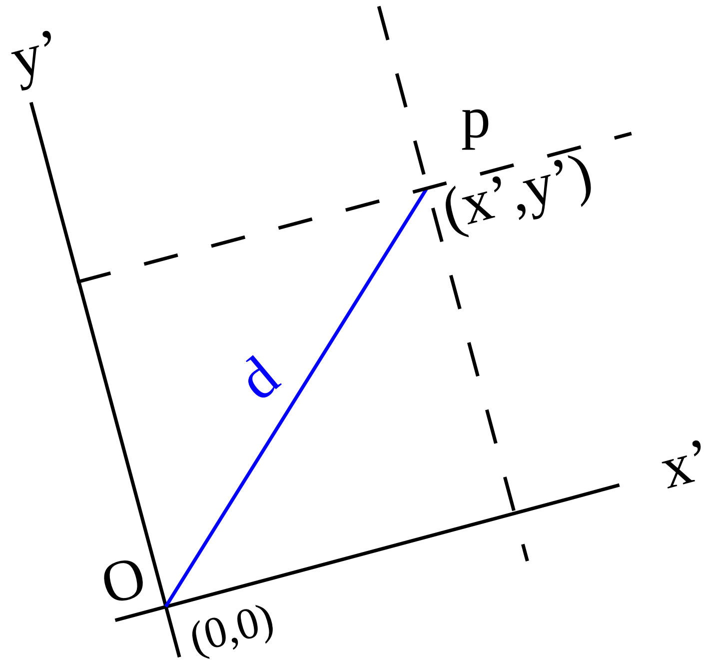
5. metric $\Leftarrow$ Einstein field equations

GR: coordinate transformations



GR: coordinate transformations



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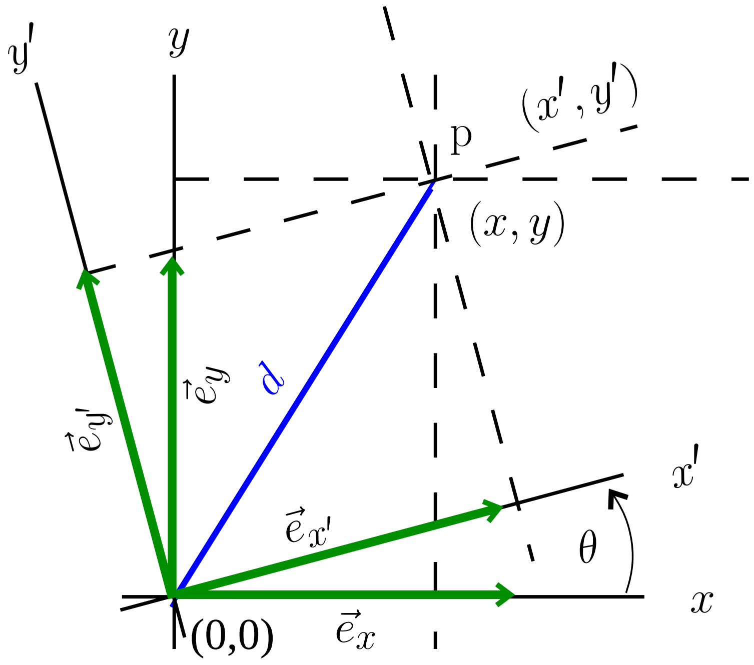

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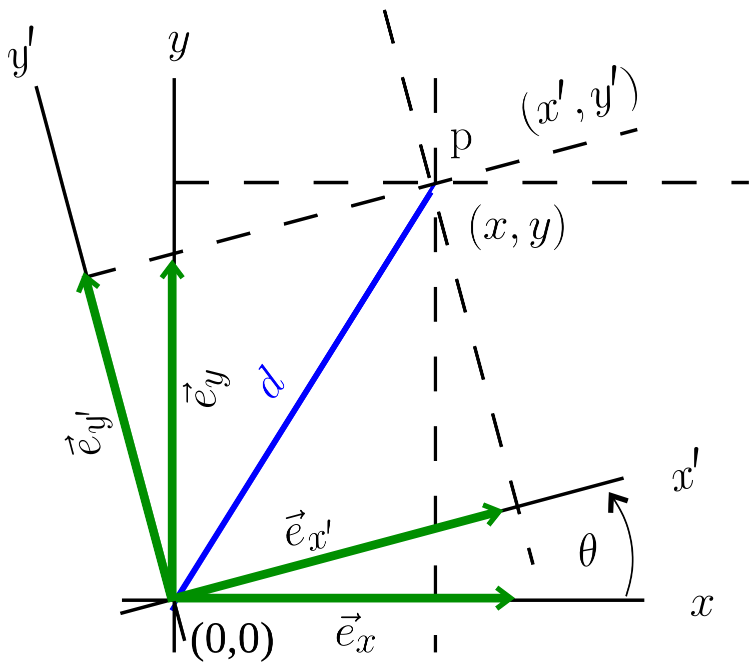

$$\begin{pmatrix} x' \\ y' \end{pmatrix} = \begin{pmatrix} \cos\theta & \sin\theta \\ -\sin\theta & \cos\theta \end{pmatrix} \begin{pmatrix} x \\ y \end{pmatrix}$$

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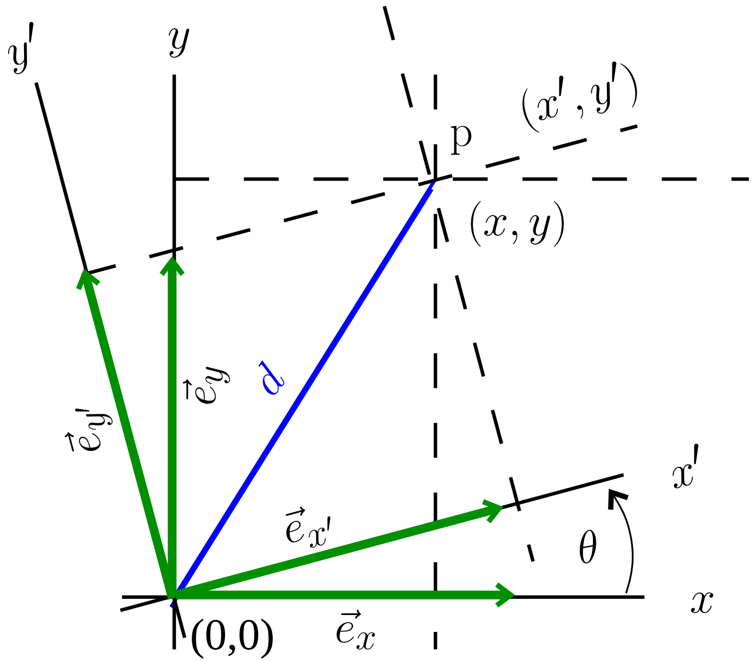

$$\begin{pmatrix} x' \\ y' \end{pmatrix} = \Lambda \begin{pmatrix} x \\ y \end{pmatrix}$$

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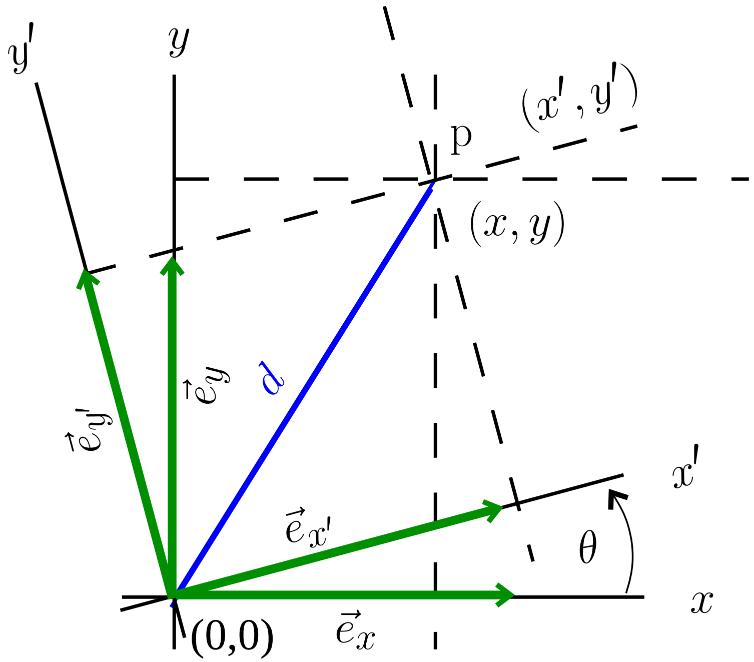
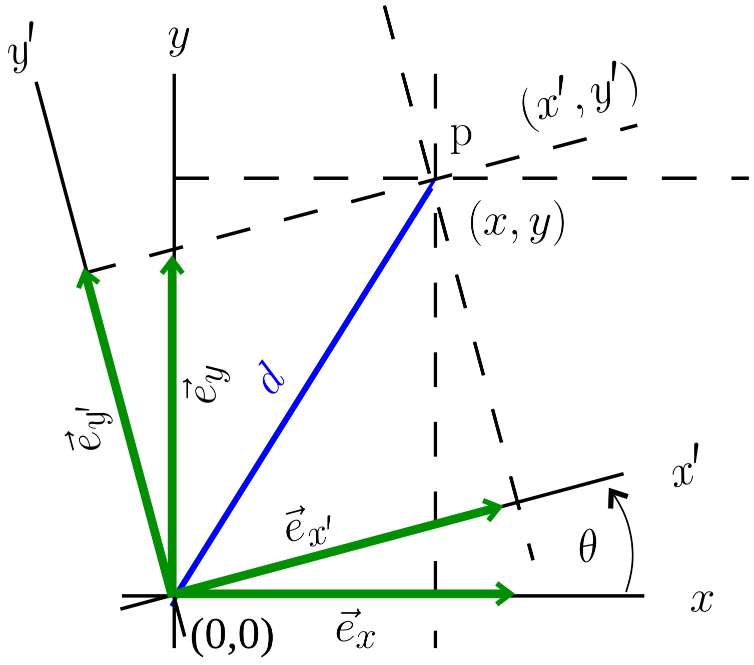


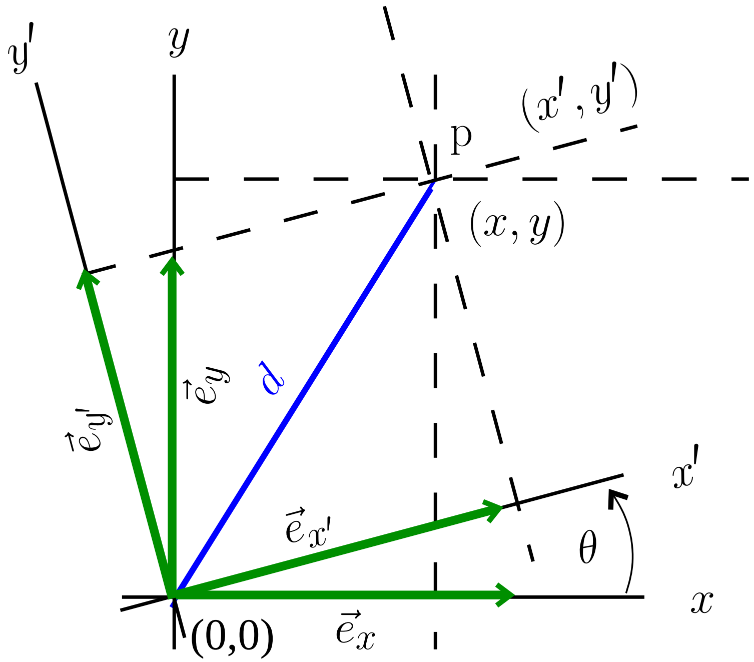
but $\begin{pmatrix} \cos\theta \\ \sin\theta \end{pmatrix} = \begin{pmatrix} \cos\theta & -\sin\theta \\ \sin\theta & \cos\theta \end{pmatrix} \begin{pmatrix} 1 \\ 0 \end{pmatrix}$

GR: coordinate transformations



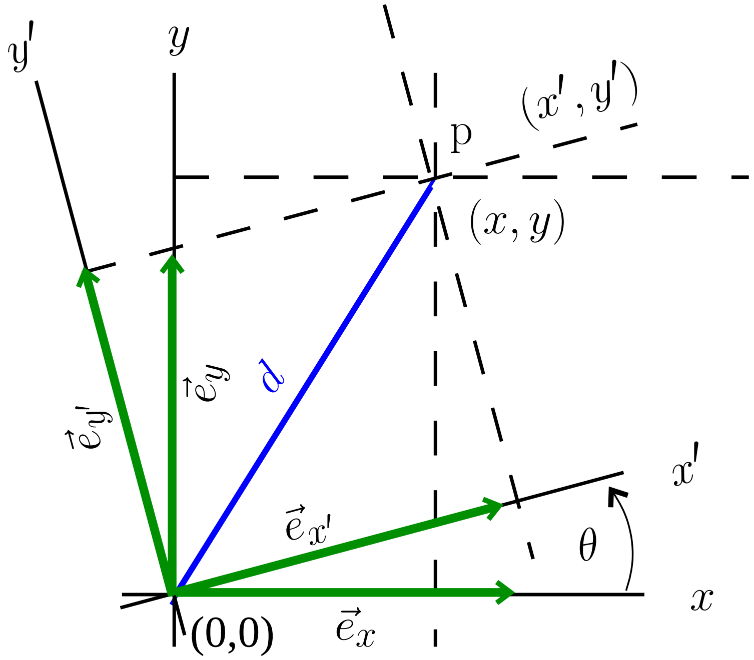
$$\begin{pmatrix} \cos\theta \\ \sin\theta \end{pmatrix} = \begin{pmatrix} \cos\theta & -\sin\theta \\ \sin\theta & \cos\theta \end{pmatrix} \begin{pmatrix} 1 \\ 0 \end{pmatrix} + \begin{pmatrix} \cos\theta & -\sin\theta \\ \sin\theta & \cos\theta \end{pmatrix} \begin{pmatrix} 0 \\ 1 \end{pmatrix}$$

GR: coordinate transformations



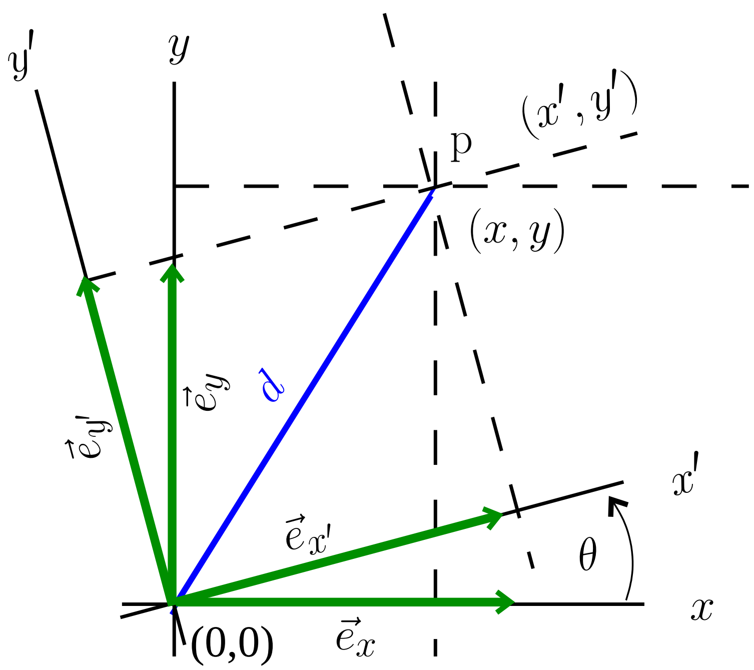
$$\vec{e}_{x'} = \begin{pmatrix} \cos\theta & -\sin\theta \\ \sin\theta & \cos\theta \end{pmatrix} \vec{e}_x + \begin{pmatrix} \cos\theta & -\sin\theta \\ \sin\theta & \cos\theta \end{pmatrix} \vec{e}_y$$

GR: coordinate transformations



$$\vec{e}_{x'} = \Lambda^{x}_{x'}\,\vec{e}_{x} + \Lambda^{y}_{x'}\,\vec{e}_{y}$$

GR: coordinate transformations



also:

$$\begin{pmatrix} -\sin\theta \\ \cos\theta \end{pmatrix} = \begin{pmatrix} \cos\theta & -\sin\theta \\ \sin\theta & \cos\theta \end{pmatrix} \begin{pmatrix} 1 \\ 0 \end{pmatrix} + \begin{pmatrix} \cos\theta & -\sin\theta \\ \sin\theta & \cos\theta \end{pmatrix} \begin{pmatrix} 0 \\ 1 \end{pmatrix}$$

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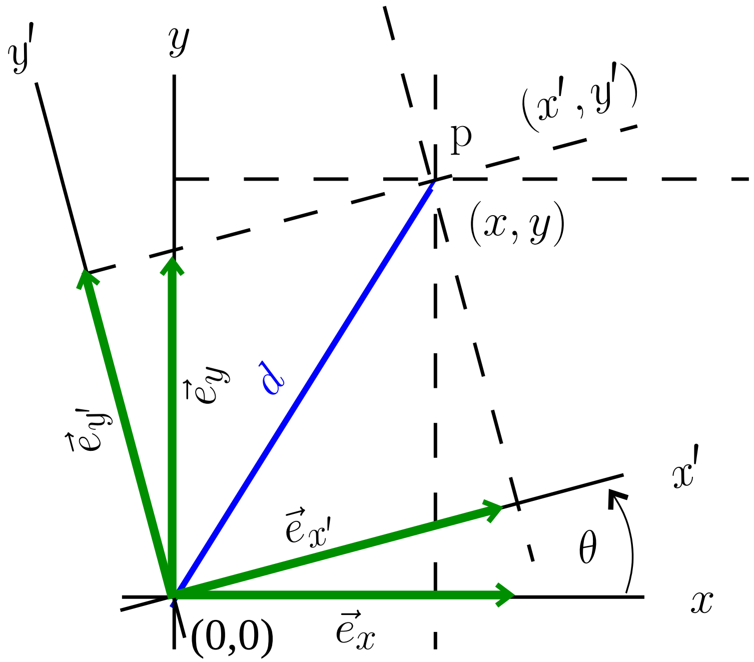

$$\vec{e}_{y'} = \begin{pmatrix} \cos\theta & -\sin\theta \\ \sin\theta & \cos\theta \end{pmatrix} \vec{e}_x + \begin{pmatrix} \cos\theta & -\sin\theta \\ \sin\theta & \cos\theta \end{pmatrix} \vec{e}_y$$

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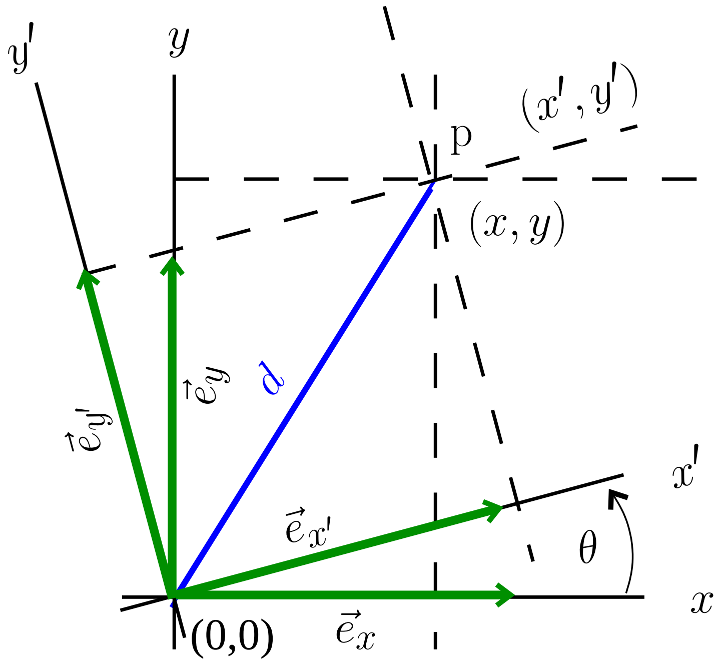

summary:

$$\vec{e}_{x'} = \Lambda^x_{x'}\,\vec{e}_x + \Lambda^y_{x'}\,\vec{e}_y,$$

$$\vec{e}_{y'} = \Lambda^x_{y'}\,\vec{e}_x + \Lambda^y_{y'}\,\vec{e}_y,$$

where $\Lambda^\alpha_{\beta'}$:= element of inverse of $\Lambda^{\alpha'}_\beta$,

$$\begin{pmatrix} x' \\ y' \end{pmatrix} = \Lambda \begin{pmatrix} x \\ y \end{pmatrix}$$

GR: coordinate transformations

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$$\vec{p} \rightarrow_{\mathcal{O}'} \begin{pmatrix} x' \\ y' \end{pmatrix} = \Lambda \begin{pmatrix} x \\ y \end{pmatrix}$$

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$$\vec{p} = \sum_i p^i \vec{e}_i$$

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$\vec{p} = p^i \vec{e}_i$ (w:Einstein summation)

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Einstein summation:

- **coordinates** like r, θ, x, y:

not a sum: $\Lambda^{x}_{y'}\,\vec{e}_x$

- repeated up-down coordinate **indices** like $i, j \in \{0, 1, 2\}$ or $\alpha, \beta, \gamma, \lambda, \mu, \nu \in \{0, 1, 2, 3\}$:

sum: $\Lambda^{i}_{j'}\,\vec{e}_i := \Lambda^{x}_{y'}\,\vec{e}_x + \Lambda^{y}_{y'}\,\vec{e}_y$ for a 2D manifold, coords x, y

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$\vec{e}_{\mu'} = \sum_\nu \Lambda^{\nu}{}_{\mu'} \vec{e}_\nu$

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$$\vec{e}_{x'} = \Lambda^x_{x'}\,\vec{e}_x + \Lambda^y_{x'}\,\vec{e}_y, \qquad \vec{e}_{y'} = \Lambda^x_{y'}\,\vec{e}_x + \Lambda^y_{y'}\,\vec{e}_y,$$

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new coords of vector $\vec{p}$ = $\Lambda\times$ old coords of same vector $\vec{p}$

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vector invariance requires contravariance of its coords
“contra" = inverse of change of basis vectors

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vector invariance requires contravariance of its coords
“contra" = inverse of change of basis vectors

- $\vec{p}$ is invariant: no dependence on coords
- $\vec{p}$ is contravariant: p^i change inversely to $\vec{e}_i$

GR: coord. transf.: 1-forms

ϕ = scalar field = $\phi(x, y) \equiv \phi(x', y')$

write $\phi_{,x} := \frac{\partial \phi}{\partial x} =: (\tilde{\mathrm{d}}\phi)_x$

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What is the relation between $(\phi_{,x'}, \phi_{,y'})$

and $(\phi_{,x}, \phi_{,y})$?

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ϕ depends either on x and y, or on x' and y'

$$\Rightarrow \frac{\partial\phi}{\partial x'} = \frac{\partial\phi}{\partial x}\frac{\partial x}{\partial x'} + \frac{\partial\phi}{\partial y}\frac{\partial y}{\partial x'}$$

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$(\phi_{,x'}, \phi_{,y'}) =$

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$(\phi_{,x'}, \phi_{,y'}) = \quad (\phi_{,x}\, x_{,x'} + \phi_{,y}\, y_{,x'} \,,\, \phi_{,x}\, x_{,y'} + \phi_{,y}\, y_{,y'})$

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$$\begin{pmatrix} x \\ y \end{pmatrix} = \begin{pmatrix} \cos\theta & -\sin\theta \\ \sin\theta & \cos\theta \end{pmatrix} \begin{pmatrix} x' \\ y' \end{pmatrix}$$ (example: rotation)

$x_{,x'} = \frac{\partial x}{\partial x'} = \cos\theta$

$x_{,y'} = \frac{\partial x}{\partial y'} = -\sin\theta \ldots$

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$$\begin{pmatrix} x \\ y \end{pmatrix} = \Lambda^{-1} \begin{pmatrix} x' \\ y' \end{pmatrix} \quad \text{(general)}$$

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$$\tilde{\mathrm{d}}\phi = \left((\tilde{\mathrm{d}}\phi)_{x'}, (\tilde{\mathrm{d}}\phi)_{y'}\right) = \left((\tilde{\mathrm{d}}\phi)_x, (\tilde{\mathrm{d}}\phi)_y\right) \Lambda^{-1}$$

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$$(\tilde{\mathrm{d}}\phi)_{\mu'} = (\tilde{\mathrm{d}}\phi)_\nu \Lambda^\nu{}_{\mu'}$$

GR: coord. transf.: 1-forms

basis vectors of different bases: $\vec{e}_{\mu'} = \Lambda^{\nu}{}_{\mu'} \vec{e}_{\nu}$

same vector: $(\vec{p})^{\mu'} = \Lambda^{\mu'}{}_{\nu} (\vec{p})^{\nu}$

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w:Covariance and contravariance of vectors

GR: $\vec{p}, \tilde{q}, \langle \vec{p}, \tilde{q} \rangle$, $\mathbf{g}$

GR tensors: two different scalar products

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can be called I with components δ^μ_ν in a coordinate basis

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think: vector $\rightarrow$ column vector
1-form $\rightarrow$ row vector

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$\langle\ ,\ \rangle$ = (1,1)-tensor = “row-column" matrix I with $I^\mu_\nu = \delta^\mu_\nu$

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GR tensors: two different scalar products

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ordinary linear algebra: column vectors, row vectors, matrices

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(m, n)-tensor algebra: m column n row $m + n$-arrays

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(m, n)-tensor algebra: m column n row $m + n$-arrays

e.g.: $(0, 2)$-tensor: metric $g_{\mu\nu}$

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using $\langle \ , \ \rangle$, $(1, 0)$-tensor = vector = function of 1-forms

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loosely speaking, the second $\otimes$ means “function of two vectors" (or 1-forms, or a vector and a 1-form) in *that particular left-to-right order*

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warning: the “rank" of tensors has two different meanings: w:Tensor_(intrinsic_definition)#Tensor_rank

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dimension of $V^* \otimes V^*$ = 16 (for $V =$ spacetime)

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also written: $\vec{A} \cdot \vec{B}$ “dot product"

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$\mathbf{g}(\vec{A}, \vec{B}) \;\; = A^r B^r + A^\theta B^\theta r^2 \;\; = A^x B^x + A^y B^y$

in general, for a 2-form **T**, $\mathbf{T}(\vec{A}, \vec{B}) \neq \mathbf{T}(\vec{B}, \vec{A})$

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$$\mathbf{g} = \sum_{\mu \in \{r,\theta\}, \nu \in \{r,\theta\}} g_{\mu\nu} \tilde{e}^\mu \otimes \tilde{e}^\nu$$

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g can be applied to basis vectors $\vec{e}_\mu$

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we can define components (used earlier): $g_{\mu\nu} := \mathbf{g}(\vec{e}_\mu, \vec{e}_\nu)$

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e.g. $\mathbf{g} = g_{rr}\tilde{e}^r \otimes \tilde{e}^r + g_{r\theta}\tilde{e}^r \otimes \tilde{e}^\theta + g_{\theta r}\tilde{e}^\theta \otimes \tilde{e}^r + g_{\theta\theta}\tilde{e}^\theta \otimes \tilde{e}^\theta$

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check: $\mathbf{g}(\vec{e}_r, \vec{e}_r) = g_{rr}$?

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lower index: $g_{\mu\nu}A^\mu = A_\nu$

raise index: $g^{\mu\nu}B_\nu = B^\mu$

GR: what is a coordinate?

a coordinate, e.g. x^0 or x^1 is a scalar field on the 4-manifold

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a coordinate system x^μ = set of four scalar fields on the 4-manifold

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(Bertschinger writes $x^\mu_{\mathbf{x}}$ to show dependence on position $\mathbf{x}$ in manifold $\neq$ vector space)

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x^μ are differentiable *almost everywhere*

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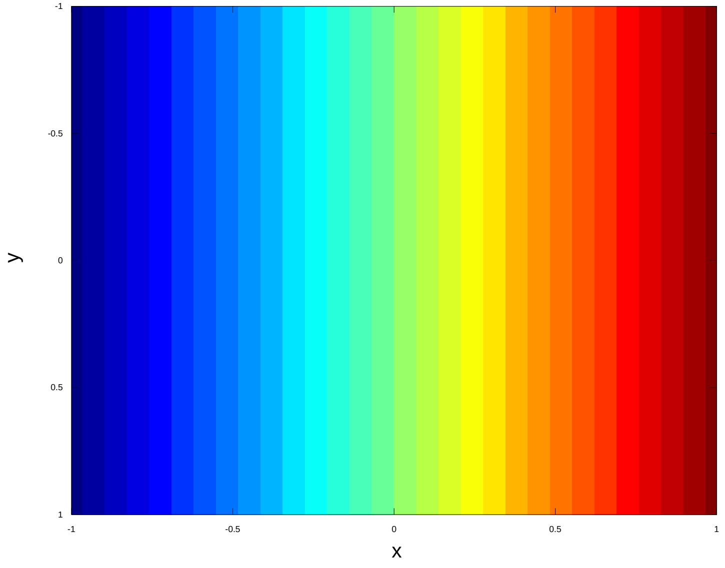

e.g. on $\mathbb{R}^2$

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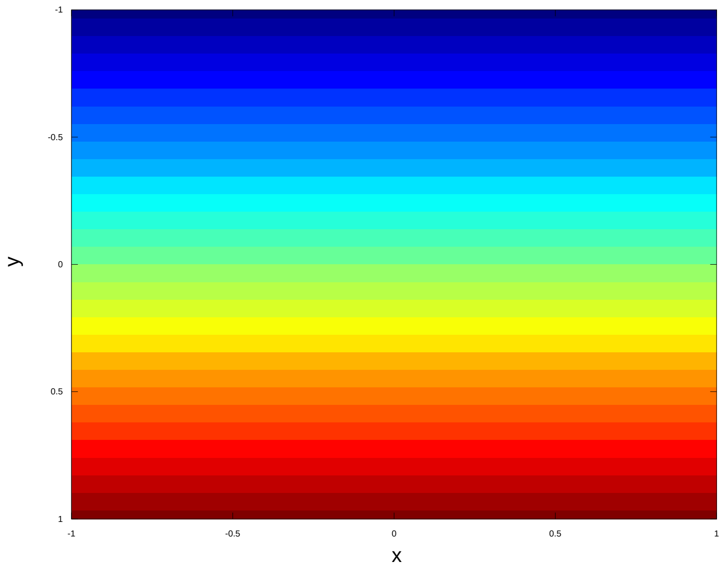

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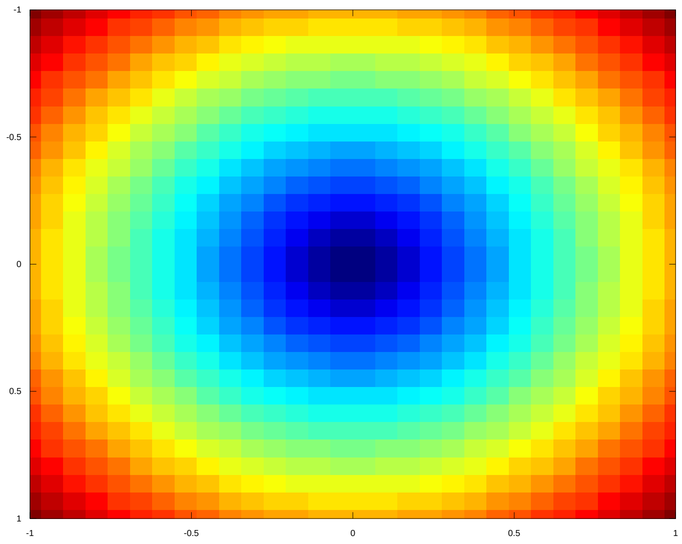

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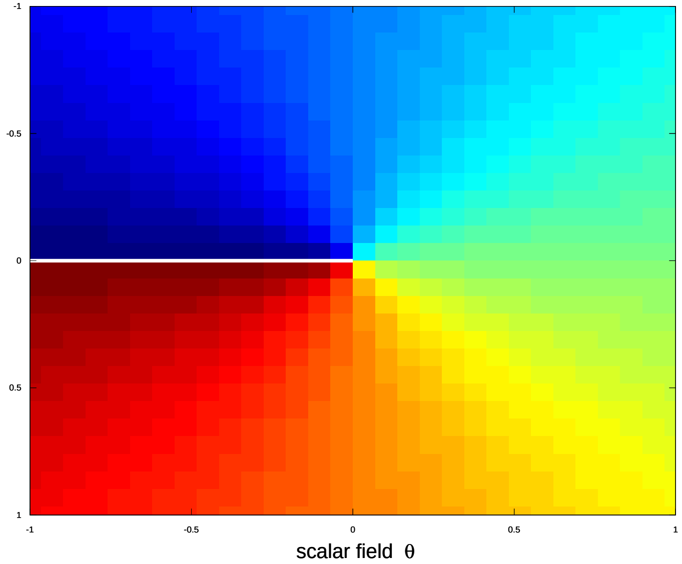

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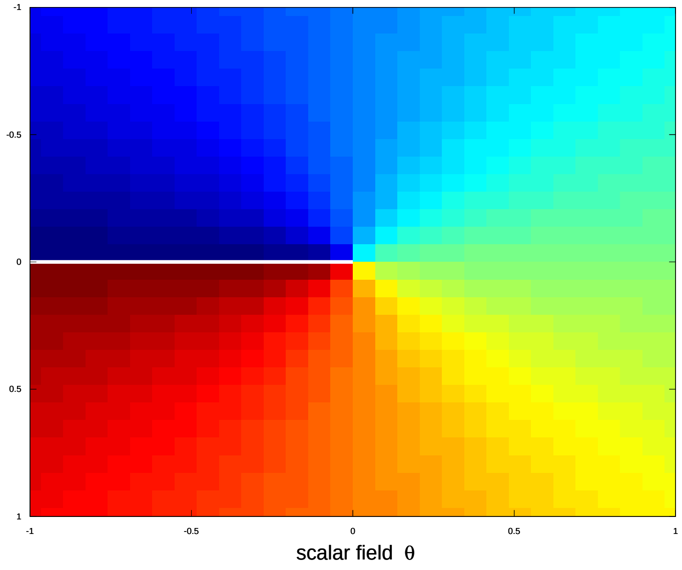

e.g. on $\mathbb{R}^2$

coordinate singularity $\neq$ singularity in manifold

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where $\tilde{\mathrm{d}} = \tilde{e}^\mu \partial_\mu$ in a coordinate basis

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(Bertschinger writes $\widetilde{\nabla}$ for the gradient $\tilde{\mathrm{d}}$)

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$\mathrm{d}s^2 = g_{\mu\nu}\mathrm{d}x^\mu \mathrm{d}x^\nu$ if x^μ are a coordinate basis

GR: e.g. Euclidean g on $\mathbb{R}^2$

$g_{r\theta}$ and g_{xy}

$$\mathrm{d}s^2 = \mathrm{d}x^2 + \mathrm{d}y^2 = \mathrm{d}r^2 + r^2 \mathrm{d}\theta^2$$

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$\mathrm{d}s^2 = \mathrm{d}x^2 + \mathrm{d}y^2 = \mathrm{d}r^2 + r^2\mathrm{d}\theta^2$

$\vec{e}_x \cdot \vec{e}_x = 1 = \vec{e}_y \cdot \vec{e}_y$, others zero

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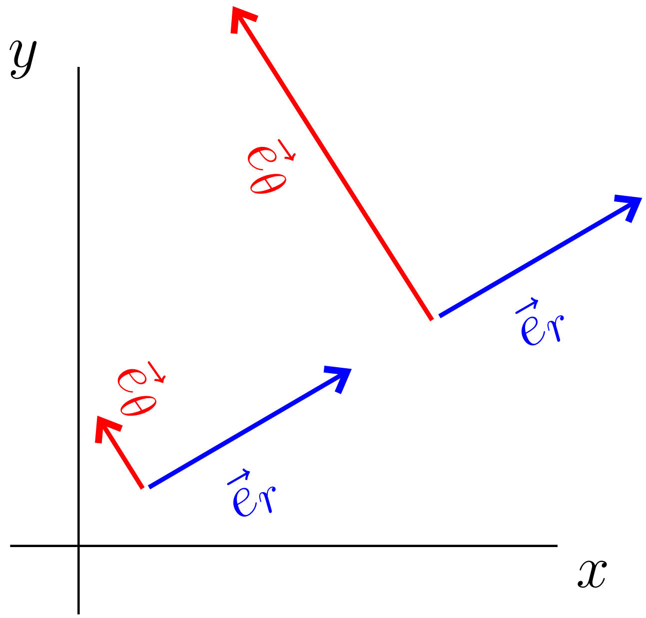
$g_{r\theta}$ and g_{xy}

$\mathrm{d}s^2 = \mathrm{d}x^2 + \mathrm{d}y^2 = \mathrm{d}r^2 + r^2\mathrm{d}\theta^2$

$\vec{e}_x \cdot \vec{e}_x = 1 = \vec{e}_y \cdot \vec{e}_y$, others zero

$\vec{e}_r \cdot \vec{e}_r = 1$, $\vec{e}_\theta \cdot \vec{e}_\theta = r^2 \neq 1$



GR: e.g. Euclidean g on $\mathbb{R}^2$

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$g^{\mu\alpha} g_{\alpha\nu} = \delta^\mu_\nu \Rightarrow g^{xx} = 1 = g^{yy}, g^{xy} = 0 = g^{yx}$

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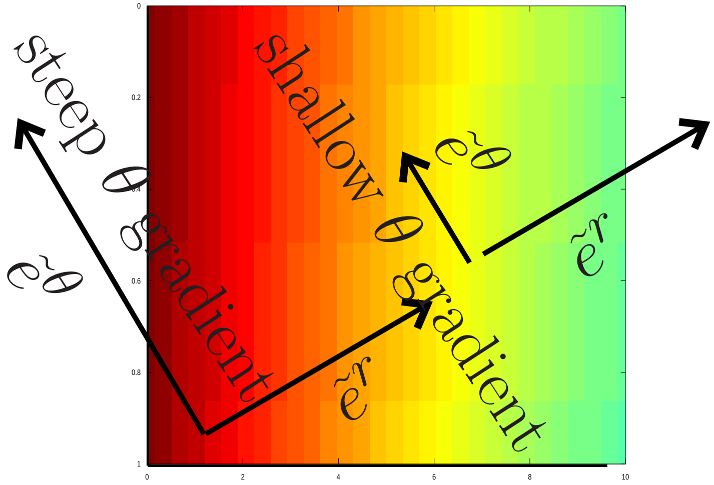

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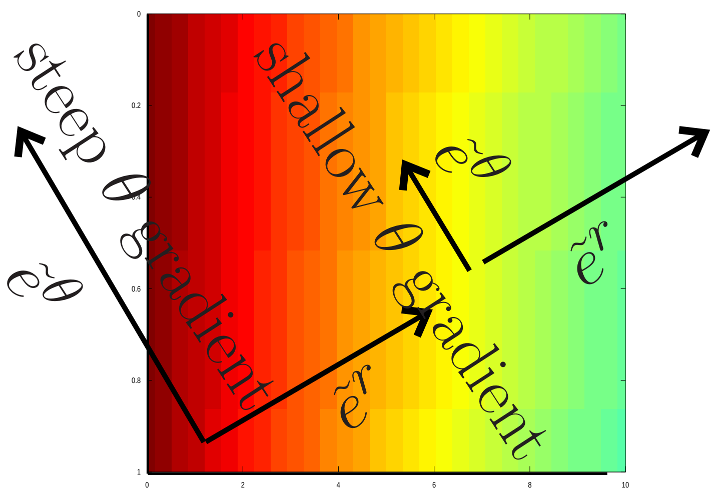

so $\tilde{e}^r \cdot \tilde{e}^r = 1$, $\tilde{e}^\theta \cdot \tilde{e}^\theta = r^{-2} \neq 1$

GR: gradient of a vector: $\nabla\vec{A}$

gradient of scalar field: $\tilde{\mathrm{d}}\phi \equiv \widetilde{\nabla}\phi$

GR: gradient of a vector: $\nabla\vec{A}$

what is gradient of vector field $\widetilde{\nabla}\vec{A}$?

GR: gradient of a vector: $\nabla \vec{A}$

$$\widetilde{\nabla} \vec{A} = \widetilde{\nabla}(A^\nu \vec{e}_\nu)$$

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$$\widetilde{\nabla}\vec{A} = \widetilde{\nabla}(A^\nu \vec{e}_\nu)$$

$$= \tilde{e}^\mu \partial_\mu (A^\nu \vec{e}_\nu)$$

$$= \tilde{e}^\mu \otimes [(\partial_\mu A^\nu)\vec{e}_\nu + A^\nu \partial_\mu \vec{e}_\nu]$$ by product rule and linearity

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$$= \partial_\mu A^\nu \tilde{e}^\mu \otimes \vec{e}_\nu + A^\nu \tilde{e}^\mu \otimes \partial_\mu \vec{e}_\nu$$

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$$= \partial_\mu A^\nu\tilde{e}^\mu \otimes \vec{e}_\nu + A^\nu\tilde{e}^\mu \otimes \partial_\mu\vec{e}_\nu$$

give a name to the second part: it must be a linear combination of basis vectors $\vec{e}_\lambda$

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define $\Gamma^\lambda{}_{\nu\mu} \vec{e}_\lambda := \partial_\mu \vec{e}_\nu$ Christoffel symbols of second kind (symmetric defn)

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$= \partial_\mu A^\nu\tilde{e}^\mu \otimes \vec{e}_\nu + A^\nu\Gamma^\lambda{}_{\nu\mu}\tilde{e}^\mu \otimes \vec{e}_\lambda$ since any $\Gamma^\lambda{}_{\nu\mu}$ is a scalar

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since name of summation index is arbitrary, e.g. $\sum_\lambda x^{-2\lambda} = \sum_\mu x^{-2\mu} = \sum_\nu x^{-2\nu}$

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$$\boxed{\nabla_\mu A^\nu := A^\nu{}_{;\mu} := \partial_\mu A^\nu + A^\lambda \Gamma^\nu{}_{\lambda\mu}}$$

w:covariant derivative of vector

GR: gradient of a vector: $\nabla\vec{A}$

$\widetilde{\nabla}\vec{A} = \widetilde{\nabla}(A^\nu \vec{e}_\nu)$

$= \tilde{e}^\mu \partial_\mu (A^\nu \vec{e}_\nu)$

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define $\Gamma^\lambda{}_{\nu\mu} \vec{e}_\lambda := \partial_\mu \vec{e}_\nu$ Christoffel symbols of second kind (symmetric defn)

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mathematically deeper: $\tilde{\nabla}$, usually written just as ∇, is the w:Levi-Civita connection

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and $\widetilde{\nabla}$ on a $(1, 0)$-tensor field = vector field $\vec{A}$ gives a $(1, 1)$-tensor with components $\nabla_{\mu}A^{\nu} = \partial_{\mu}A^{\nu} + A^{\lambda}{\Gamma^{\nu}}_{\lambda\mu}$

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- not components of tensor: $\Gamma^{\nu}{}_{\lambda\mu}$

GR: gradient of one-form $\widetilde{\nabla}\tilde{A}$

how does a one-form change with position? $\widetilde{\nabla}\tilde{A} =$?

GR: gradient of one-form $\widetilde{\nabla}\tilde{A}$

evaluating $\widetilde{\nabla}\tilde{A}$ as we did $\widetilde{\nabla}\vec{A}$ shows that we again need $\partial_\mu \tilde{e}^\nu = F^\nu{}_{\lambda\mu}\tilde{e}^\lambda$ for some coefficients $F^\nu{}_{\lambda\mu}$

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relation between vectors and one-forms: $\langle \tilde{e}^\nu, \vec{e}_\lambda \rangle = \delta^\nu{}_\lambda$

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$\partial_\mu \delta^\nu{}_\lambda = 0$ (obviously)

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can we use the product rule with this scalar product?

$\partial_\mu \left(\left\langle \tilde{A}, \vec{B} \right\rangle \right) =$?

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can we use the product rule with this scalar product?

$\partial_\mu \left(\left\langle \tilde{A}, \vec{B} \right\rangle \right) = \partial_\mu \left(A_\nu B^\nu \right)$ in some coordinate basis

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$$\partial_\mu \left(\left\langle \tilde{A}, \vec{B} \right\rangle \right) = \partial_\mu \left(A_\nu B^\nu \right)$$

$= (\partial_\mu A_\nu) B^\nu + A_\nu (\partial_\mu B^\nu)$ by product rule on functions

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$= \left\langle \partial_\mu \tilde{A}, \vec{B} \right\rangle + \left\langle \tilde{A}, \partial_\mu \vec{B} \right\rangle$ since

$\partial_\mu \tilde{A} = (\partial_\mu A_0, \partial_\mu A_1, \partial_\mu A_2, \partial_\mu A_3)$

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GR: gradient of one-form $\widetilde{\nabla}\tilde{A}$

evaluating $\widetilde{\nabla}\tilde{A}$ as we did $\widetilde{\nabla}\vec{A}$ shows that we again need $\partial_\mu \tilde{e}^\nu = F^\nu{}_{\lambda\mu}\tilde{e}^\lambda$ for some coefficients $F^\nu{}_{\lambda\mu}$

how can we relate $\Gamma^\nu{}_{\lambda\mu}$ to $F^\nu{}_{\lambda\mu}$?

relation between vectors and one-forms: $\langle \tilde{e}^\nu, \vec{e}_\lambda \rangle = \delta^\nu{}_\lambda$

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$= \left\langle F^\nu{}_{\kappa\mu}\tilde{e}^\kappa, \vec{e}_\lambda \right\rangle + \left\langle \tilde{e}^\nu, \Gamma^\kappa{}_{\lambda\mu}\vec{e}_\kappa \right\rangle$

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$$\boxed{\nabla_\mu A^\nu = \partial_\mu A^\nu + A^\lambda \Gamma^\nu{}_{\lambda\mu} \quad , \quad \nabla_\mu A_\nu = \partial_\mu A_\nu - A_\lambda \Gamma^\lambda{}_{\mu\nu}}$$

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$$A^\nu{}_{;\mu} = A^\nu{}_{,\mu} + A^\lambda \Gamma^\nu{}_{\lambda\mu} \quad , \quad A_{\nu;\mu} = A_{\nu,\mu} - A_\lambda \Gamma^\lambda{}_{\mu\nu}$$

GR: smooth manifold and $\widetilde{\nabla}\mathbf{g}$

similarly, we can write the $(0,3)$-tensor

$$\widetilde{\nabla}\mathbf{g} = (\nabla_\lambda g_{\mu\nu})\tilde{e}^\lambda \otimes \tilde{e}^\mu \otimes \tilde{e}^\nu$$

giving $\nabla_\lambda g_{\mu\nu} = \partial_\lambda g_{\mu\nu} - \Gamma^\kappa{}_{\mu\lambda} g_{\kappa\nu} - \Gamma^\kappa{}_{\nu\lambda} g_{\mu\kappa}$

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also $\widetilde{\nabla}\mathbf{g}^{-1} = (\nabla_\lambda g^{\mu\nu})\tilde{e}^\lambda \otimes \vec{e}_\mu \otimes \vec{e}_\nu$

and $\nabla_\lambda g^{\mu\nu} = \partial_\lambda g^{\mu\nu} + \Gamma^\mu{}_{\kappa\lambda} g_{\kappa\nu} + \Gamma^\nu{}_{\kappa\lambda} g_{\mu\kappa}$

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Do we know anything interesting about $\widetilde{\nabla}\mathbf{g}$ for the manifolds of interest to GR?

GR: smooth manifold and $\widetilde{\nabla}\mathbf{g}$

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Do we know anything interesting about $\widetilde{\nabla}\mathbf{g}$ for the manifolds of interest to GR?

First, we need a rough description of the manifolds we need for GR.

GR: smooth manifold and $\widetilde{\nabla}$g

topological manifold M
w:Manifold#Mathematical_definition

- only topological properties needed

GR: smooth manifold and $\widetilde{\nabla \mathbf{g}}$

topological manifold M
w:Manifold#Mathematical_definition

- only topological properties needed
- no differentiability, no metric needed

GR: smooth manifold and $\widetilde{\nabla}\mathbf{g}$

topological manifold M

w:Manifold#Mathematical_definition

- only topological properties needed

next: relation with $\mathbb{R}^4$ (or M^4)

GR: smooth manifold and $\widetilde{\nabla}\mathbf{g}$

topological manifold M
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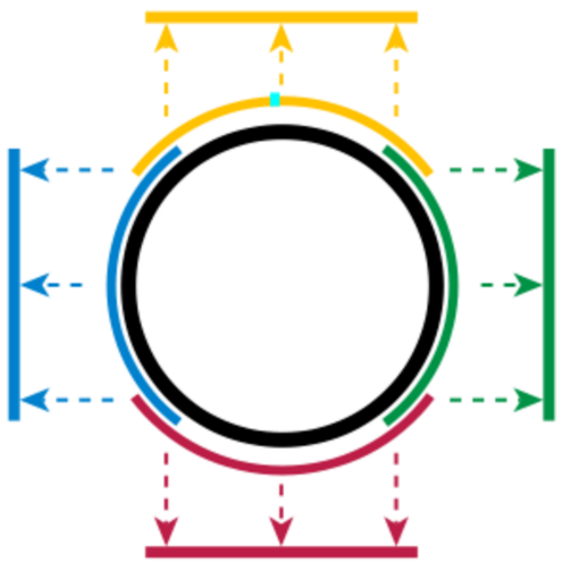
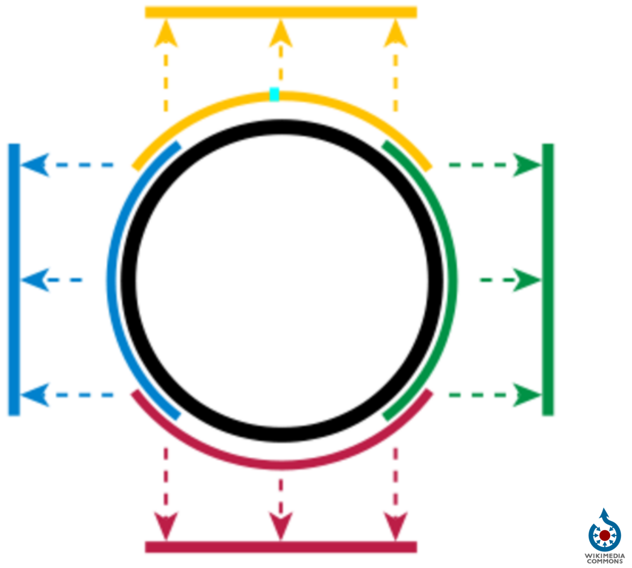


GR: smooth manifold and $\widetilde{\nabla}\mathbf{g}$

topological manifold M
w:Manifold#Mathematical_definition

- only topological properties needed

next: relation with $\mathbb{R}^4$ (or M^4)



w:Manifold

- chart := function ϕ_α from part of pseudo-4-manifold M to part of M^4 (Minkowski)
- atlas := set of overlapping charts that cover M

GR: smooth manifold and $\widetilde{\nabla}$g

if every transition chart $:= \phi_\beta \circ \phi_\alpha^{-1}$ in an atlas for M is differentiable on $\mathbb{R}^4$ (or M^4), then M is a w:differentiable 4-(pseudo-)manifold

GR: smooth manifold and $\widetilde{\nabla}\mathbf{g}$

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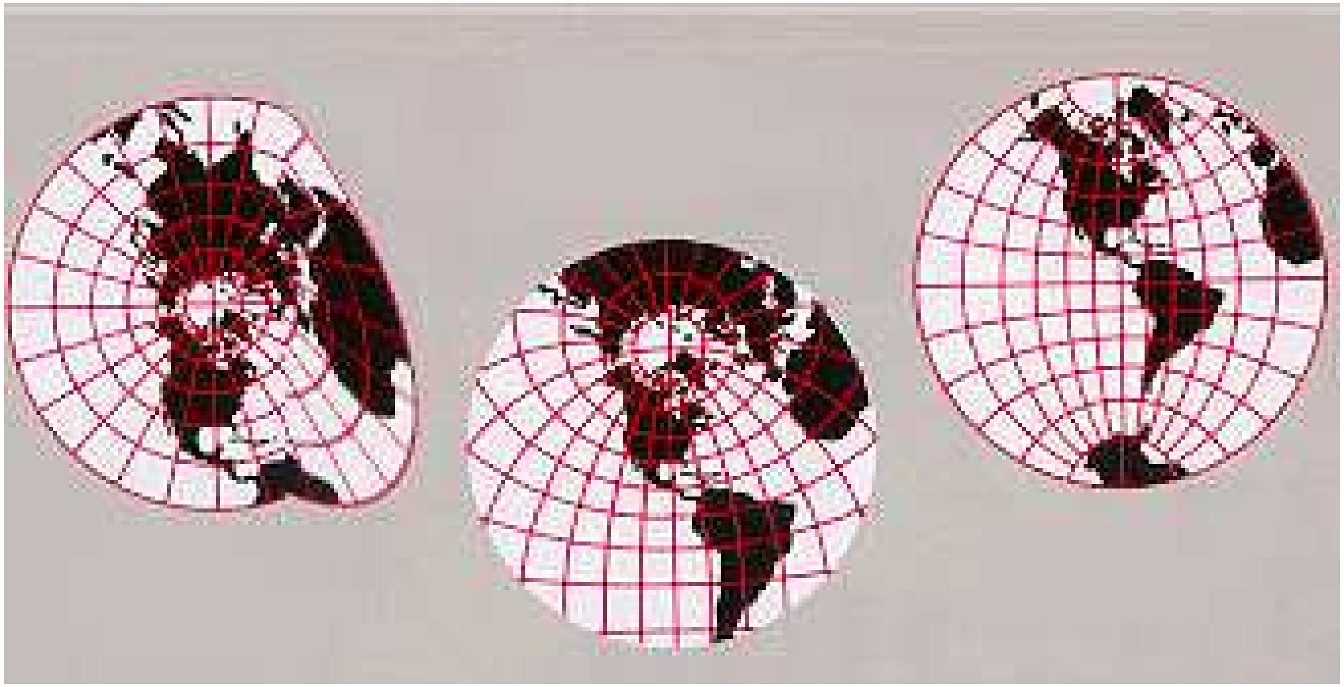

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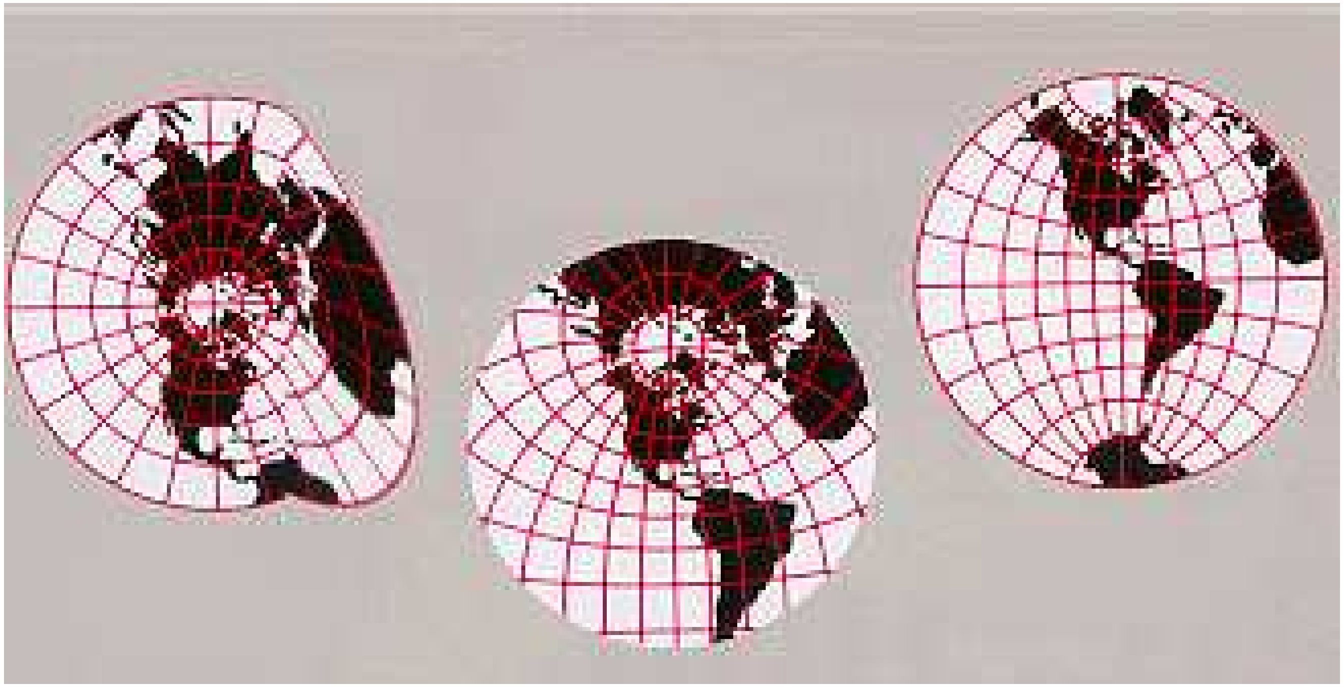

w:

projections (left-to-right) ϕ_1, ϕ_2, ϕ_3 from S^2 to $\mathbb{R}^2$

GR: smooth manifold and $\widetilde{\nabla}$g

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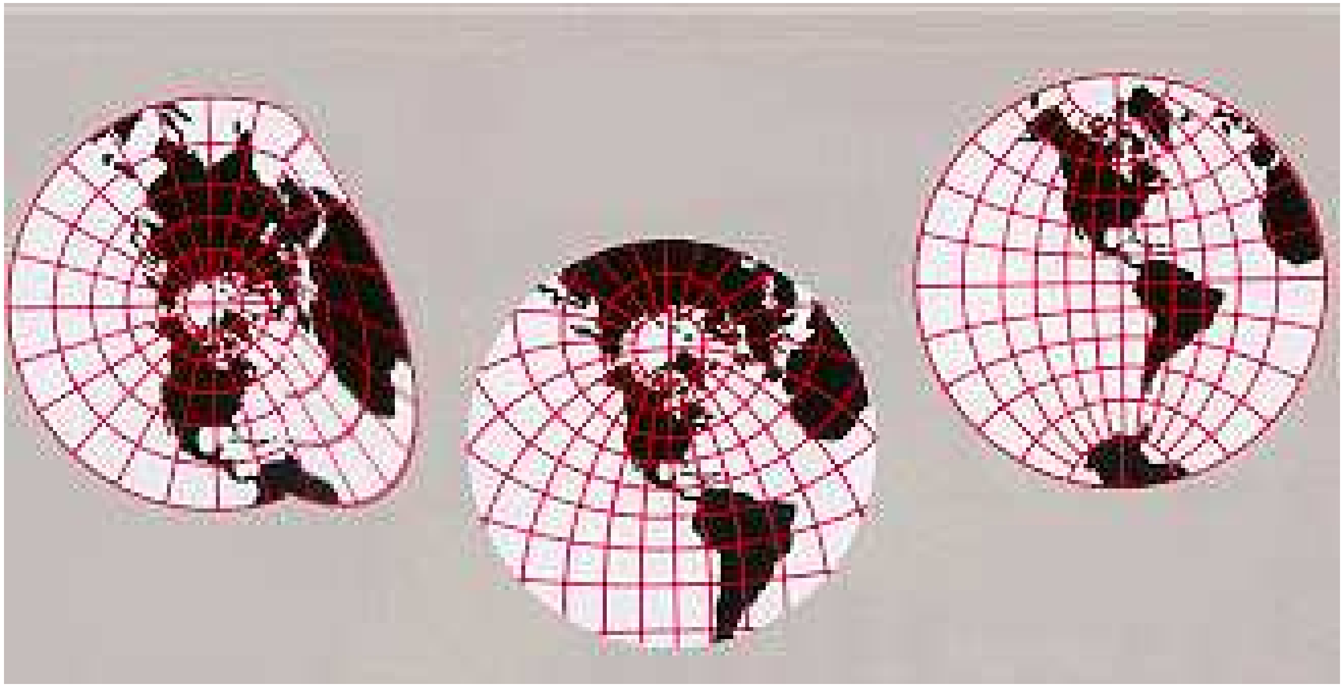

w:

ϕ_1 is not differentiable, so $\phi_1 \circ \phi_2^{-1}$ is not differentiable

GR: smooth manifold and $\widetilde{\nabla}\mathbf{g}$

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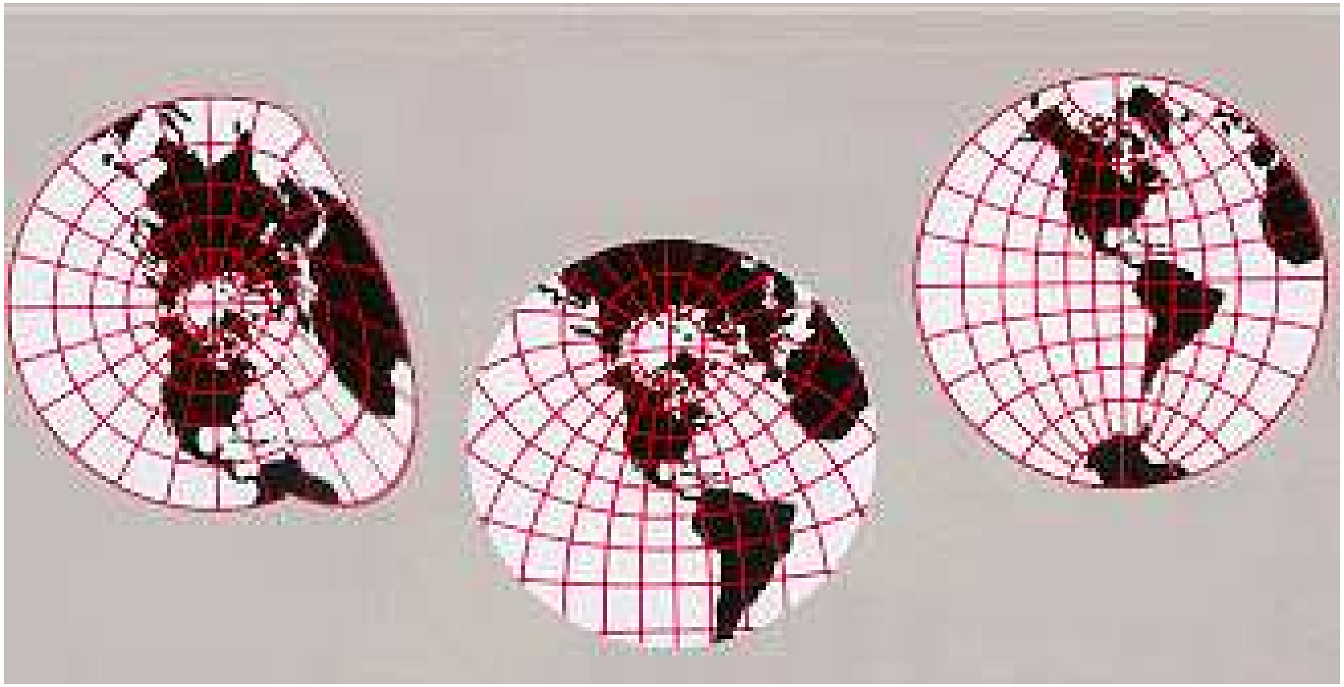

w:

atlas not enough to show that $S^2 =$ differentiable 2-manifold

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if $\forall k \geq 1$, $\exists k$-th derivatives, then M is a smooth 4-(pseudo-)manifold

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if a (pseudo-)w:Riemannian metric **g** can be added to M, then $(M,$**g**$)$ is a (pseudo-)Riemannian 4-manifold

GR: smooth manifold and $\widetilde{\nabla}$g

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if a (pseudo-)w:Riemannian metric **g** can be added to M, then $(M,\mathbf{g})$ is a (pseudo-)Riemannian 4-manifold

if **g** has signature $(1, n-1)$ (i.e. $(-,+,+,+)$ or $(+,-,-,-)$, etc.), then $(M,\mathbf{g})$ is a Lorentzian n-manifold

GR: smooth manifold and $\widetilde{\nabla}$g

topological manifolds

GR: smooth manifold and $\widetilde{\nabla}\mathbf{g}$

topological manifolds

differentiable (pseudo-)manifolds

GR: smooth manifold and $\widetilde{\nabla}\mathbf{g}$

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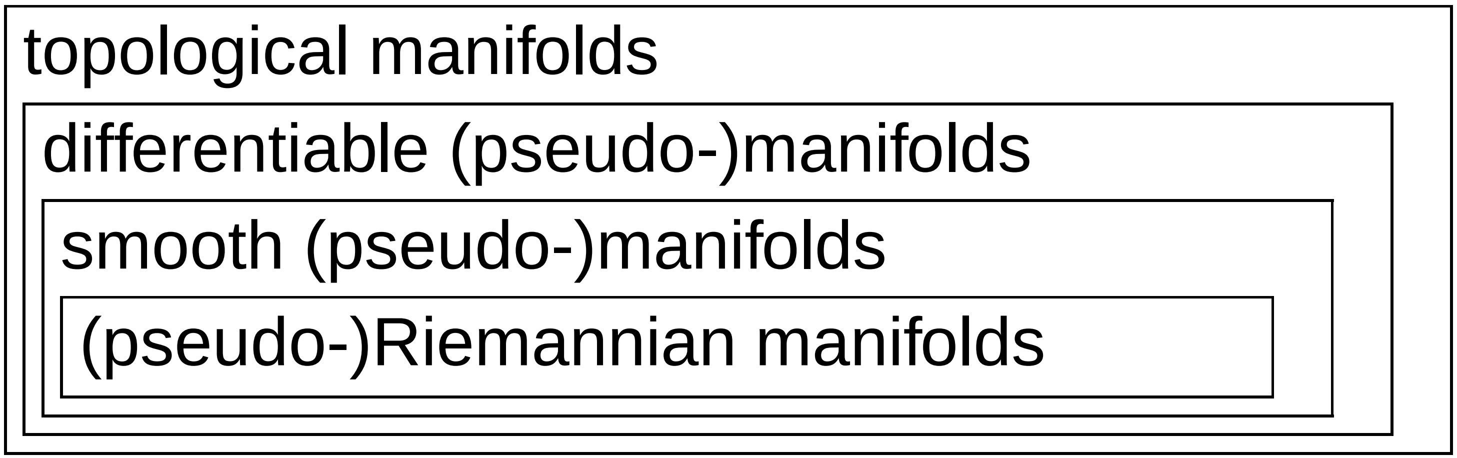
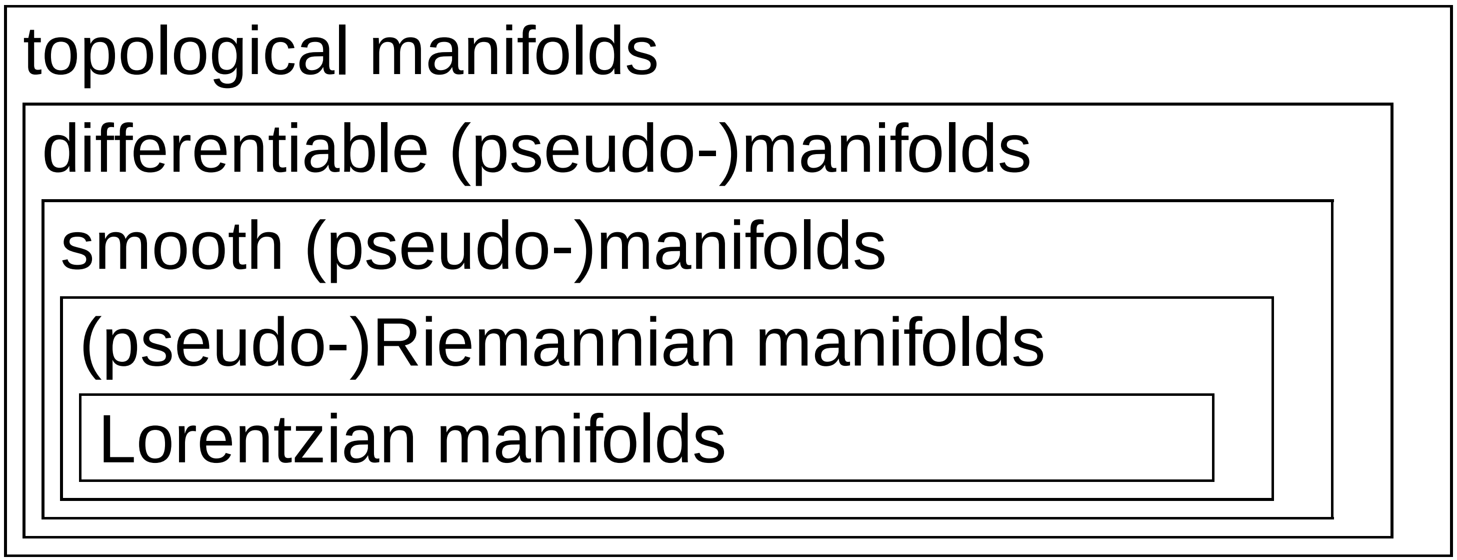
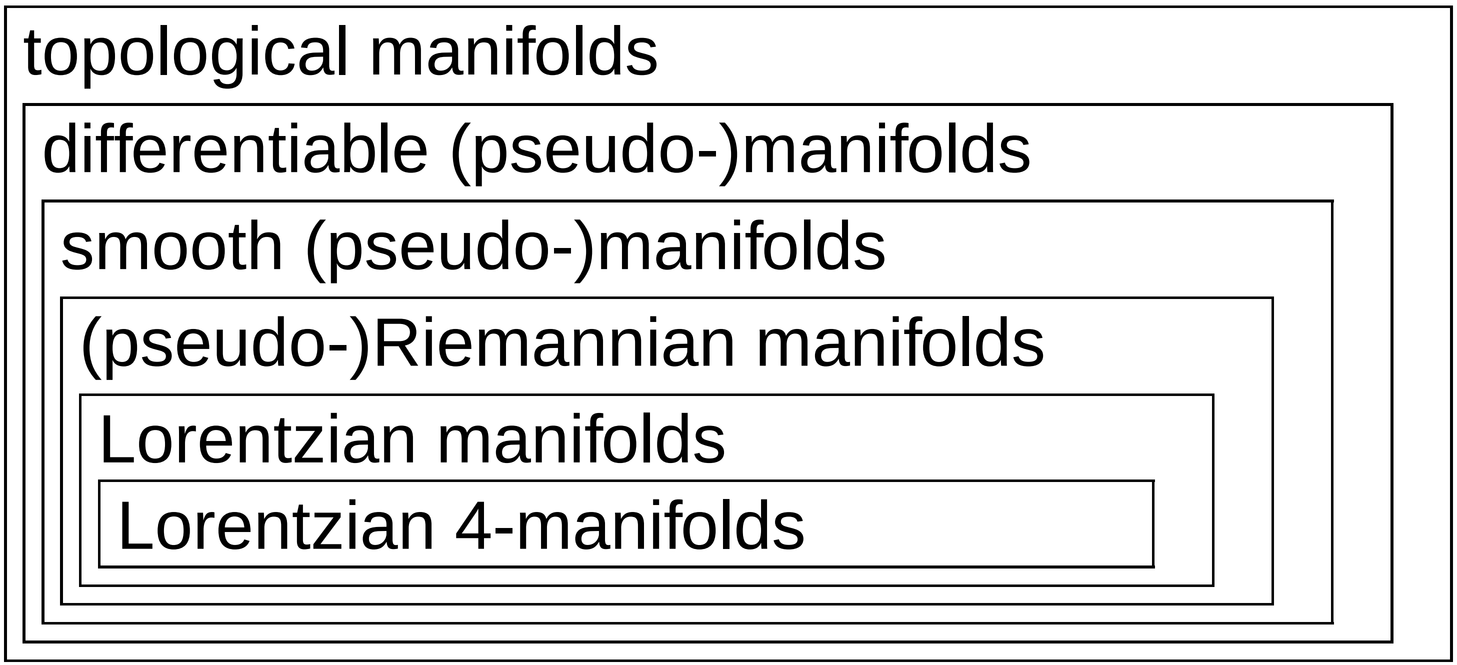


GR: smooth manifold and $\widetilde{\nabla}\mathbf{g}$



GR: smooth manifold and $\widetilde{\nabla}\mathbf{g}$



GR: smooth manifold and $\widetilde{\nabla}\mathbf{g}$

- topological manifolds
  - differentiable (pseudo-)manifolds
    - smooth (pseudo-)manifolds
      - (pseudo-)Riemannian manifolds
        - Lorentzian manifolds
          - Lorentzian 4-manifolds

GR: assume that spacetime is a Lorentzian 4-manifold

GR: smooth manifold and $\widetilde{\nabla}\mathbf{g}$

from above:

$$\nabla_\lambda g_{\mu\nu} = \partial_\lambda g_{\mu\nu} - \Gamma^\kappa{}_{\mu\lambda} g_{\kappa\nu} - \Gamma^\kappa{}_{\nu\lambda} g_{\mu\kappa}$$

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in the tangent space at $\mathbf{x}$, $\exists$ coordinate basis $\vec{e}_{\bar{\mu}}$ with

$$g_{\bar{\mu}\bar{\nu}} = \eta_{\bar{\mu}\bar{\nu}} = \mathrm{diag}(-1,1,1,1) = g^{\bar{\mu}\bar{\nu}}$$

$$\Rightarrow \partial_{\bar{\lambda}} g_{\bar{\mu}\bar{\nu}} = \partial_{\bar{\lambda}} \eta_{\bar{\mu}\bar{\nu}}$$

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also, $\Gamma^{\bar{\lambda}}{}_{\bar{\nu}\bar{\mu}} \vec{e}_{\bar{\lambda}} := \vec{e}_{\bar{\nu},\bar{\mu}} = \partial_{\bar{\mu}} \vec{e}_{\bar{\nu}}$

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but in a Cartesian or Minkowski (vector) space, the basis vectors always point in the same direction and their lengths are fixed

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M$^4 \Rightarrow \Gamma^{\bar{\lambda}}{}_{\bar{\nu}\bar{\mu}} \vec{e}_{\bar{\lambda}} = 0$

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so $\nabla_{\bar{\lambda}} g_{\bar{\mu}\bar{\nu}} = 0$

so $\widetilde{\nabla}\mathbf{g} = \mathbf{0}$ (also $\widetilde{\nabla}\mathbf{g}^{-1} = 0$) on the tangent space, since if true in one coord system, also true in others

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$\Rightarrow \partial_{\bar{\lambda}} g_{\bar{\mu}\bar{\nu}} = \partial_{\bar{\lambda}} \eta_{\bar{\mu}\bar{\nu}} \;= 0$

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so $\nabla_{\bar{\lambda}} g_{\bar{\mu}\bar{\nu}} = 0$

so $\widetilde{\nabla}\mathbf{g} = \mathbf{0} = \widetilde{\nabla}\mathbf{g}^{-1}$ on tangent space

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GR: smooth manifold and $\widetilde{\nabla}\mathbf{g}$

from above:

$\nabla_\lambda g_{\mu\nu} = \partial_\lambda g_{\mu\nu} - \Gamma^\kappa{}_{\mu\lambda} g_{\kappa\nu} - \Gamma^\kappa{}_{\nu\lambda} g_{\mu\kappa}$

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...

$\Gamma^\lambda{}_{\mu\nu} = \frac{1}{2} g^{\lambda\kappa} (\partial_\mu g_{\nu\kappa} + \partial_\nu g_{\mu\kappa} - \partial_\kappa g_{\mu\nu})$ in a coordinate basis

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warning: $\{x^\mu(\lambda)\}$ at some λ on the manifold is a point on the manifold but NOT a vector; while $\mathrm{d}\vec{x}$ — in the tangent space — IS a vector

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using $\vec{V}(\lambda) := \frac{\mathrm{d}\vec{x}}{\mathrm{d}\lambda}$, project covariant derivative to curve
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∇_V written by Bertschinger without $^{\rightarrow}$or $^{\sim}$ because ∇_V **T** of tensor **T** has the same tensor order as **T**

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so in a coord basis,

$$\boxed{\nabla_V \vec{A} = \left(\frac{\mathrm{d}A^\nu}{\mathrm{d}\lambda} + V^\mu A^\kappa {\Gamma^\nu}_{\kappa\mu} \right) \vec{e}_\nu}$$

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special (interesting) case: vector field $\vec{A}$ and curve with tangents $\vec{V} := \frac{\mathrm{d}\vec{x}}{\mathrm{d}\lambda}$ where $\vec{A}$ “locally does not change direction"

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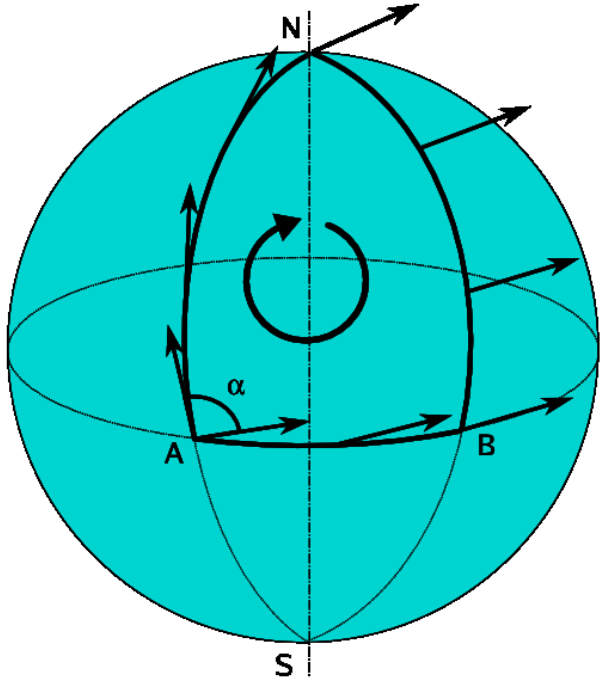
i.e. $\nabla_V \vec{A} = 0$

$\nabla_V \vec{A} = 0$ defn: parallel transport of $\vec{A}$ along path $\mathbf{x}(\lambda)$

where $\vec{V} := \frac{\mathrm{d}\vec{x}}{\mathrm{d}\lambda}$

GR: directional deriv.: $< \widetilde{\nabla} \vec{A}, \vec{V} >$

example:



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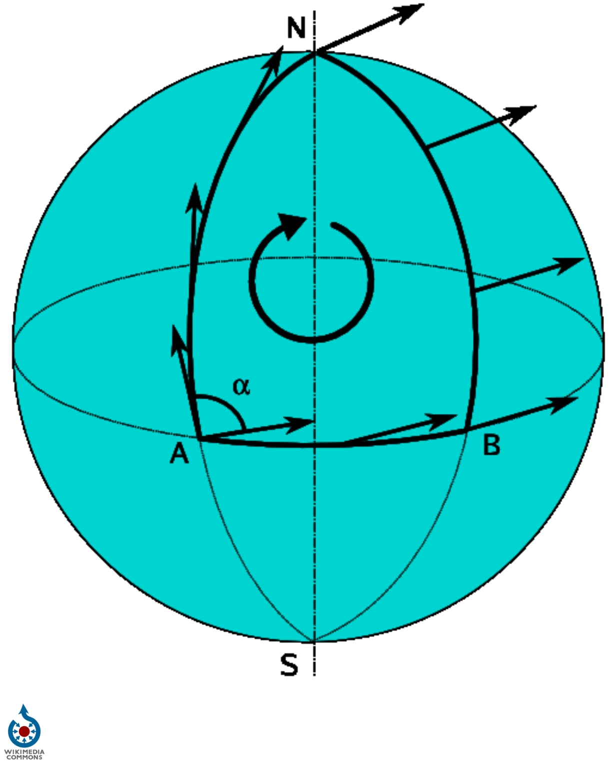

on S^2, parallel transport of $\vec{A}$ around a closed loop does not conserve $\vec{A}$'s direction

GR: directional deriv.: $< \widetilde{\nabla} \vec{A}, \vec{V} >$

$\boxed{\nabla_V \vec{V} = 0}$ defn: w:geodesic

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- more general definition of “straight line" than “shortest distance between two points"
- tensorial definition — independent of coordinate basis
- allows more than one “straight line" between two points a and b in a manifold — consider S^2, T^3

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a and b in a manifold — consider S^2, T^3

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i.e. $\frac{\mathrm{d}^2 x^\nu}{\mathrm{d}\lambda^2} + \frac{\mathrm{d}x^\mu}{\mathrm{d}\lambda}\frac{\mathrm{d}x^\kappa}{\mathrm{d}\lambda}\Gamma^\nu{}_{\kappa\mu} = 0 \; \forall \nu$

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cf w:Euler-Lagrange equation

GR: parallel transp. @closed curve

parallel transport around “small" parallelogram in two directions $\mathrm{d}\vec{x}_1$, $\mathrm{d}\vec{x}_2$,

(“1" and “2" are not component indices here)

GR: parallel transp. @closed curve

parallel transport around “small" parallelogram in two directions $\mathrm{d}\vec{x}_1$, $\mathrm{d}\vec{x}_2$,

What is the change in $\vec{A}$ after parallel transport around the closed loop $\mathrm{d}\vec{x}_1$, $\mathrm{d}\vec{x}_2$, $-\mathrm{d}\vec{x}_1$, $-\mathrm{d}\vec{x}_2$?

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$\propto \mathrm{d}\vec{x}_1, \mathrm{d}\vec{x}_2$

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$\Rightarrow$ must exist a tensor **R** that is a function of 3 vectors (“inputs"),

i.e. is a $\otimes$ of 3 one-forms

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defn: $\boxed{-\mathbf{R}(\cdot, \vec{A}, \mathrm{d}\vec{x}_1, \mathrm{d}\vec{x}_2) := \mathrm{d}\vec{A}(\cdot)}$

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(minus sign convention: MTW1973, Bertschinger)

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$$\mathrm{d}\vec{A}(\cdot) = -R^{\mu}_{\ \ \nu\alpha\beta} A^{\nu}\, \mathrm{d}x_1^{\alpha}\, \mathrm{d}x_2^{\beta}\, \vec{e}_{\mu}(\cdot)$$

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i.e. $-\mathbf{R} = R^{\mu}{}_{\nu\alpha\beta} \vec{e}_{\mu} \otimes \tilde{e}^{\nu} \otimes \tilde{e}^{\alpha} \otimes \tilde{e}^{\beta}$

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GR: Riemann tensor

how can **R** be evaluated?

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use covariant derivatives of covariant derivatives . . .

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Ricci identity:

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- so **R** has second order partial derivatives of $g_{\nu\kappa}, \ldots$

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- second order ∂:

(pseudo-)manifold globally like $\mathbb{R}^3$ (M^4) $\Leftrightarrow R^{\mu}{}_{\nu\alpha\beta}(\mathbf{x}) = 0 \ \forall \mathbf{x}$

GR: Bianchi, Ricci, Einstein

. . . second Bianchi identity:

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warning: “R" written coordinate-free (without indices) may mean:

- an order 4, dimension 256 tensor **R**;
- an order 2, dimension 16 tensor **R** or R; or
- an order 0, dimension 1 tensor $\equiv$ scalar R
- all three are fields over a spacetime 4-manifold

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defn Einstein tensor (by components):

$$\boxed{G^{\mu\nu} := R^{\mu\nu} - \frac{1}{2} g^{\mu\nu} R}$$

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Cactus - http://cactuscode.org